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JAN 78 F S CHWIEROTH, G L ZARUR, A NAGAL N00173-77-C-0008

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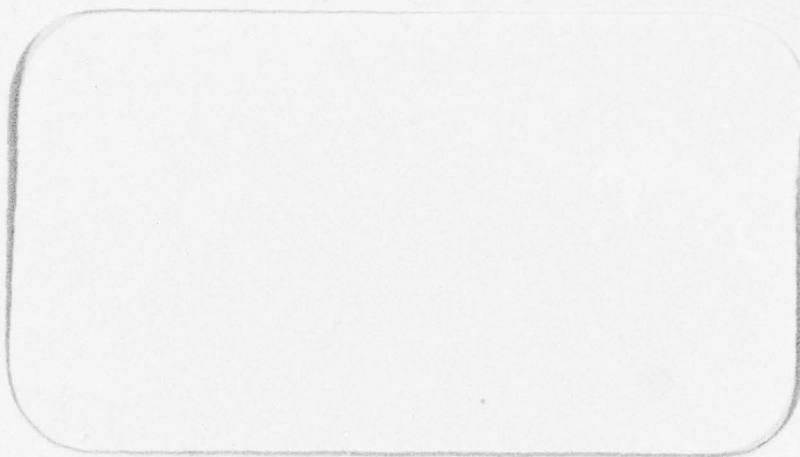
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⑥ A FORTRAN Code for the  
Calculation of Sound Propagation  
in a Range Dependent Ocean, II,  
Range Functions and Mode Coupling.

Volume

(First Version)

⑫ 57 p.

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20. ABSTRACT (Continue on reverse side if necessary and identify by block number) Pierce's adiabatic normal-mode theory of sound propagation in a range dependent ocean duct is extended to the case of sound channels with less than gradual range dependence. A computational method is devised to solve the ensuing coupled range equations by diagonalization. This is applied to a channel with arbitrary (numerically given) range dependence, by dividing it into range segments with constant properties each. The local depth functions are taken as the Airy function solutions in a piece-wise linear-		



ized stratified medium with arbitrary sound velocity profile. The method is applied to some simple as well as to realistic cases of range dependent sound channels, and the dependence of the mode coupling effects on the degree of range dependence is determined.

Appended to the theory part exhibited in this present first version of the report is a listing of a computer program based on the foregoing theory; it computes range functions and solves the coupled-mode program of sound propagation in a range dependent ocean.

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Mode Coupling in a Range Dependent  
Under-Ocean Sound Channel

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Abstract

✓ Pierce's adiabatic normal-mode theory of sound propagation in a range dependent ocean duct is extended to the case of a sound channel with less than gradual range dependence. A computational method is devised to solve the ensuing coupled range equations by diagonalization. This is applied to a channel with arbitrary (numerically given) range dependence, by dividing it into range segments with constant properties each. The local depth functions are taken as the Airy function solutions in a piece-wise linearized stratified medium with arbitrary sound velocity profile. The method is applied to some simple as well as to realistic cases of range dependent sound channels, and the dependence of the mode coupling effects on the degree of range dependence is determined. ↵

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## Introduction

Normal mode theory of underwater sound propagation can, in its original form<sup>1</sup>, be applied only to situations where both the ocean floor and the sound velocity profile do not depend on the horizontal coordinate (range coordinate), so that the wave equation may be separated into modal range functions and depth functions. In a realistic ocean, this is very rarely the case, but an approximate method has been developed by Pierce<sup>2</sup> in which gradually range dependent environments can be handled to first approximation. This "adiabatic" approximation, based on the Born-Oppenheimer method of molecular physics, utilizes a modal expansion in which the depth functions are taken as the "local" modes of a locally stratified ocean which depend parametrically on range. Insertion into the wave equation then leads to a set of coupled equations for the range functions, but these become uncoupled if in the adiabatic approximation, the coupling terms are neglected. The adiabatic method has been applied in our previous work to a wedge-shaped continental shelf with constant sound velocity<sup>3</sup>, to a sound channel with parabolic profile that opens up linearly in range<sup>4</sup>, and to sound channels with arbitrary (but gradual) dependence of the ocean floor and of the sound speed profile (which may depend arbitrarily on depth also) on the range<sup>5</sup>.

No energy transfer between modes takes place in the adiabatic approximation. If the mode coupling terms are not neglected, the range equations are coupled and energy flows between the modes.



This has been discussed on the basis of coupled power equations by McDaniel<sup>6</sup>, and was considered by her in a calculation of statistical mode conversion effects as caused by a rough ocean floor<sup>7</sup>. For the deterministic case, the mode coupling problem has been solved in our earlier work<sup>8</sup> on the sound channel with a range dependent parabolic velocity profile, by using diagonalization techniques for the decoupling of the range equations.

In the present work, similar techniques are used for the more general case where the profile depends arbitrarily on both depth and range, and the ocean floor on range. This is achieved by dividing the ocean into horizontal layers as well as range segments. Linearization is employed in the depth layers, and has been applied also to the range dependent quantities in the range segments in our earlier work<sup>5</sup> on the adiabatic case. The present paper, however, employs the approximation of range segment-wise constancy of the depth functions and their eigenvalues as well as of the coupling terms, in order not to encumber the more involved decoupling problem even further. A different approach to the mode coupling problem based on the "slab method"<sup>9</sup>, is due to Kanabis<sup>10</sup>.

Pierce's method will be implemented for the coupled-mode problem of an arbitrarily range dependent sound channel in the following, and applied to simple as well as to realistic cases in order to determine the dependence of the mode coupling effects on the degree of range dependence.



# I. Local Eigenmodes of the Wave Equation

We seek a solution for the time harmonic pressure field  $p(\vec{r})$  of a point source, located at the position  $\vec{r}_0$  in an inhomogeneous ocean of density  $\rho(\vec{r})$  and propagation constant

$$k(\vec{r}) = \omega / c(\vec{r}) \quad , \text{ as determined by the wave equation}$$

$$\rho(\vec{r}) \vec{\nabla} \cdot [\rho^{-1}(\vec{r}) \vec{\nabla} p(\vec{r})] + k^2(\vec{r}) p(\vec{r}) = \delta(\vec{r} - \vec{r}_0). \quad (1a)$$

Adopting a cylindrical coordinate system with the z-axis pointing vertically downward, we shall assume the sound speed  $c(\vec{r})$  to depend arbitrarily on both the depth  $z$ , and on the horizontal range coordinate  $\vec{\rho} \equiv (x, y)$ . If the medium (which includes the ocean bottom) is layered, the boundaries of the  $\ell$ th layer  $z = z_{\ell}^{\pm}(\vec{\rho})$  may also depend on  $\vec{\rho}$  while the ocean surface  $z = z_s$  is flat. Densities  $\rho(z)$  are taken constant within each layer; they were chosen equal to  $\rho_w$  in each water layer but could be set equal to a different constant in each bottom layer. Defining a velocity potential  $\phi(\vec{r})$  such that

$$\vec{v}(\vec{r}) = - \vec{\nabla} \phi(\vec{r}),$$

$$p(\vec{r}) = \rho(\vec{r}) \partial \phi(\vec{r}) / \partial t \quad (1b)$$

( $\vec{v}$  being the particle velocity), one has in each layer a Helmholtz equation

$$\nabla^2 \phi(\vec{r}) + k^2(z, \vec{\rho}) \phi(\vec{r}) = \delta(\vec{r} - \vec{r}_0). \quad (1c)$$

Following Pierce<sup>2</sup>, we attempt a solution in the form of a quasi-separated normal-mode sum

$$\phi(\vec{r}) = \sum_n \psi_n(\vec{\rho}) u_n(z, \vec{\rho}), \quad (2)$$

in which the "local depth functions"  $u_n(z, \vec{\rho})$  satisfy at each range point  $\vec{\rho}$  (considered as a parameter) the depth equations pertaining to a stratified medium,

$$\frac{\partial^2 u_n(z, \vec{\rho})}{\partial z^2} + K_n^2(z, \vec{\rho}) u_n(z, \vec{\rho}) = 0. \quad (3)$$

The eigenvalues  $k_n(\vec{\rho})$  of the local modes enter through the local vertical wave number

$$K_n(z, \vec{\rho}) = [k^2(z, \vec{\rho}) - k_n^2(\vec{\rho})]^{1/2} \quad (4)$$

of the  $n$ th mode. They are determined by the boundary conditions satisfied by the local depth functions as used in Eq. (2),

$$\rho_\ell u_n(z_\ell^+, \vec{\rho}) = \rho_{\ell+1} u_n(z_{\ell+1}^-, \vec{\rho}), \quad (5)$$

$$[\partial u_n / \partial z]_{z_\ell^+} = [\partial u_n / \partial z]_{z_{\ell+1}^-},$$

at the boundary between layers  $\ell$  and  $\ell+1$  ( $z_\ell^+(\vec{\rho})$  being the boundary as approached from the  $\ell$ th layer, and  $z_{\ell+1}^-(\vec{\rho})$  as approached from the  $\ell+1$ st layer), corresponding to continuity of pressure and normal particle velocity. Inserting Eq.(2) into Eq. (1c) and using the orthogonality of the depth functions which follows from Eqs. (5),

$$\int_{z_s}^{\infty} [\rho(z)/\rho_w] u_n(z, \vec{\rho}) u_m(z, \vec{\rho}) dz = \delta_{nm}, \quad (6)$$

we obtain, following our earlier derivation<sup>5</sup>, the coupled system of (partial differential) range equations

$$[\nabla_{\vec{\rho}}^2 + k_n^2(\vec{\rho})] \psi_n(\vec{\rho}) = \delta(\vec{\rho}) [\rho(z_0)/\rho_w] u_n(z_0, 0) \quad (7)$$

$$-2 \sum_m [\vec{\nabla}_{\vec{\rho}} \psi_m(\vec{\rho})] \cdot \vec{M}'_{nm}(\vec{\rho}) - \sum_m \psi_m(\vec{\rho}) M''_{nm}(\vec{\rho}),$$

in which the coupling coefficients

$$\vec{M}'_{nm}(\vec{\rho}) = \int_{z_s}^{\infty} [\rho(z)/\rho_w] u_n(z, \vec{\rho}) \vec{\nabla}_{\vec{\rho}} u_m(z, \vec{\rho}) dz \quad (8a)$$

$$M''_{nm}(\vec{\rho}) = \int_{z_s}^{\infty} [\rho(z)/\rho_w] u_n(z, \vec{\rho}) \nabla_{\vec{\rho}}^2 u_m(z, \vec{\rho}) dz \quad (8b)$$

couple the  $n$ th and the  $m$ th modes. We here used the notation

$\vec{\nabla}_{\vec{\rho}} \equiv (\partial/\partial x, \partial/\partial y, 0)$ , and took the source position as  $\vec{r}_0 = (z_0,$

$\vec{\rho}_0 \equiv 0$ ). The derivation of Eq. (7), due to our use of Eqs. (5),

holds also for the present case where sloping layer boundaries are admitted. However, Eqs. (5) will guarantee the boundary conditions of the total acoustic field,

$$\rho_l \phi^{(l)}(z_l^+, \vec{\rho}) = \rho_{l+1} \phi^{(l+1)}(z_{l+1}^-, \vec{\rho}), \quad (9)$$

$$[\partial \phi^{(l)} / \partial n]_{z_l^+} = [\partial \phi^{(l+1)} / \partial n]_{z_{l+1}^-},$$

to be approximately satisfied only for gradually sloping boundaries and/or small relative density differences between layers,<sup>11</sup> i.e.,

$$\left[ (\rho_{\ell+1} - \rho_{\ell}) / \rho_{\ell+1} \right] \tan \alpha \ll 1 \quad (10)$$

(where  $\alpha$  is the angle between the boundary normal and the z-axis), as can easily be shown. This will be assumed in the following. Note, however, that no restrictions are imposed on the range gradient of the velocity profile.

The coupling terms of Eqs. (8) may be recast in a simpler form<sup>3</sup> which shows their origin from the range dependence of the sound velocity profile and of the boundaries, respectively, by employing partial integration and using Eq. (3), with the following results. For the coupling terms which are small of first order in the range dependence, one has

$$\vec{M}'_{nm}(\vec{\rho}) = \vec{C}_{nm}(\vec{\rho}) + \vec{C}_{nm}^{(s)}(\vec{\rho}) \quad (m \neq n) \quad (11a)$$

$$\vec{M}'_{nn}(\vec{\rho}) = \vec{C}_{nn}^{(s)}(\vec{\rho}) \quad (11b)$$

with the volume contribution

$$\vec{C}_{nm}(\vec{\rho}) = \vec{M}_{nm}(\vec{\rho}) / [k_m^2(\vec{\rho}) - k_n^2(\vec{\rho})], \quad (m \neq n) \quad (12a)$$



$$\vec{M}_{nm}(\vec{\rho}) = \int_{z_s}^{\infty} [\rho(z)/\rho_w] u_n(z, \vec{\rho}) [\vec{\nabla}_\rho k^2(z, \vec{\rho})] u_m(z, \vec{\rho}) dz, \quad (12b)$$

and the surface contributions

$$\tau_{nm}^{(s)}(\vec{\rho}) = M_{nm}^{(s)}(\vec{\rho}) / [k_m^2(\vec{\rho}) - k_n^2(\vec{\rho})], \quad (m \neq n) \quad (13a)$$

$$\vec{M}_{nm}^{(s)}(\vec{\rho}) = \sum_{l=1}^L (\rho_l/\rho_w) \left[ u_n \partial \vec{U}_m / \partial z - \vec{U}_m \partial u_n / \partial z \right]_{z_l^-(\vec{\rho})}^{z_l^+(\vec{\rho})} \quad (13b)$$

where L is the number of layers, and

$$\vec{U}_m(z, \vec{\rho}) = \vec{\nabla}_\rho u_m(z, \vec{\rho}); \quad (13c)$$

further,

$$\tau_{nn}^{(s)}(\vec{\rho}) = -\frac{1}{2} \sum_{l=1}^L (\rho_l/\rho_w) \left[ u_n^2(z, \vec{\rho}) \vec{\nabla}_\rho z \right]_{z_l^-(\vec{\rho})}^{z_l^+(\vec{\rho})}. \quad (13d)$$

For the second-order coupling coefficients, we shall only quote the results for the volume contributions:

$$M_{nm}''(\vec{\rho}) = N_{nm}''(\vec{\rho}) + \vec{\nabla}_\rho \cdot \vec{M}_{nm}'(\vec{\rho}), \quad (14a)$$

$$N_{nm}''(\vec{\rho}) = - \sum_{p \neq n, m} \vec{\tau}_{pn}(\vec{\rho}) \cdot \vec{\tau}_{pm}(\vec{\rho}). \quad (14b)$$

The surface contributions are absent if the boundaries are flat.

## II. Solution of the Coupled Range Equations

In our previous work on mode propagation in a range dependent environment<sup>3-5,8</sup>, two important special cases were considered in which range dependence occurs either in the Cartesian x-direction only (x-case), or in the  $\vec{\rho}$ -direction of cylindrical coordinates only ( $\rho$ -case). For sufficient distances from the source, the results of these two assumptions were shown to coincide<sup>4</sup>. We shall here restrict ourselves to the  $\rho$ -case only, so that  $\psi_n(\vec{\rho}) \equiv \psi_n(\rho)$ , and introduce the new range function  $f_n(\rho)$  via

$$\psi_n(\rho) = \rho^{-1/2} f_n(\rho). \quad (15)$$

Inserting in Eq. (7) leads to the coupled system of ordinary differential equations

$$\begin{aligned} f_n''(\rho) + (2\rho)^{-2} f_n + k_n^2(\rho) f_n = \\ = [\rho(z_0)/\rho_w] (2\pi\rho^{1/2})^{-1} d(\rho) u_n(z_0, \rho) - 2 \sum_m M'_{nm} f'_m - \sum_m V_{nm} f_m \end{aligned} \quad (16a)$$

where  $M'_{nm}(\rho)$  is defined similarly to  $\vec{M}'_{nm}$  in Eqs. (11)-(13) but with  $\vec{\nabla}_\rho$  replaced by  $d/d\rho$ , and where the further coupling terms are

$$V_{nm}(\rho) = M''_{nm}(\rho) - \rho^{-1} M'_{nm}(\rho). \quad (16b)$$

Note that as shown before<sup>4</sup>, the term  $(2\rho)^{-2} f_n$  in Eq. (16a) is negligible for distances from the source in excess of  $\sim 1$  km, for frequencies up to several Hz. Similarly, the second term in Eq. (16b) is negligible as compared with the second-last term in Eq. (16a)

in an even wider region. We write Eq. (16a) in matrix notation as

$$\left[ \frac{d^2}{d\rho^2} + (2\rho)^{-2} + k^2(\rho) \right] \vec{f}(\rho) = \vec{U} - 2M' \vec{f}' - V \vec{f} \quad (16c)$$

where  $\vec{f}(\rho)$  is a column vector with elements  $f_n(\rho)$ , and  $k^2(\rho)$  a diagonal matrix with diagonal elements  $k_n^2(\rho)$ ;  $\vec{U}$  is the column vector of the source, with components

$$U_n = \left[ \rho(z_0)/\rho_w \right] (2\pi\rho^{1/2})^{-1} \int_0^\rho u_n(z, 0). \quad (16d)$$

Our goal will be to devise a transformation which uncouples the system of Eq. (16c), at least in an approximate manner. This will be accomplished by first dividing our range  $\rho$  into a finite number of intervals with constant vertical boundaries, across which  $\phi(z)\phi$  and  $\partial\phi/\partial\rho$  have to be continuous. In the approximation of Eq. (10), this is shown by integrating over  $z$  and using orthogonality of  $u_n$ , to lead to continuity of  $f_n$  and  $df_n/d\rho$ , i.e. to

$$f_n(\rho_i^+) = f_n(\rho_{i+1}^-) \quad (17)$$

$$\left[ df_n(\rho)/d\rho \right]_{\rho_i^+} = \left[ df_n(\rho)/d\rho \right]_{\rho_{i+1}^-}$$

where  $\rho_i^+$  is the outer boundary of the  $i$ th range interval, coinciding with the inner boundary  $\rho_{i+1}^-$  of the  $i+1$ st range interval.

In our earlier study<sup>5</sup> of the "adiabatic" solution of Eq. (17) where the coupling terms were neglected, we approximated the quantity  $k_n^2(\rho) + 1/4\rho^2$  in each range by the linear expressions  $C_n\rho + D$ , joined continuously at each segment boundary where the actual values were adopted, except in the first (source) interval and in the last one (reaching to infinity), where the constant values  $k_n^2(0)$  and  $k_n^2(\infty)$ , respectively, were adopted for  $k_n^2(\rho)$ . This leads to Airy function solutions in the general intervals, and to Hankel functions  $H_0^{(1)}$  in the two extreme intervals, which become trigonometric solutions if the term  $1/4\rho^2$  is neglected. This procedure would also be possible in the present case; the coupling terms would then have the form, e.g.,

$$C_{nm}(\rho) = M_{nm} / [(C_m - C_n)\rho + D_m - D_n], \quad (18)$$

with  $M_{nm}$  depending on  $\rho$  only through  $u_n$  and  $u_m$ , since  $k^2(z, \rho)$  is also being approximated by linear expressions  $C(z)\rho + D(z)$  in order to represent the experimental range (and depth) dependent sound speed profile which is given to us only at a finite number of range points (chosen as our range interval boundaries). We shall, however, for reasons of greater simplicity approximate in each range interval the eigenvalues  $k_n^2(\rho)$  and depth functions  $u_n(z, \rho)$ , and hence also the coupling terms, by their values at the midpoints of each range interval, insuring accuracy by introducing a finer subdivision into range intervals if necessary.



We are then able to uncouple Eq. (16c) by matrix transformations involving constant transformation matrices in each range interval. Following our previous approach<sup>8</sup> used in the treatment of mode coupling for a range dependent parabolic profile, we first set

$$\vec{f}(\rho) = S_1 \vec{\Gamma}(\rho) \quad (19)$$

which, with a constant matrix  $S_1$ , leads to

$$[d^2/d\rho^2 + K^2(\rho)] \vec{\Gamma} = S_1^{-1} \vec{U} - 2\lambda \vec{\Gamma}' - W \vec{\Gamma} \quad (20)$$

where we neglected  $1/4\rho^2$ , and called

$$K^2(\rho) = S_1^{-1} k^2(\rho) S_1 \quad (21a)$$

$$\lambda = S_1^{-1} M' S_1 \quad (21b)$$

$$W = S_1^{-1} V S_1. \quad (21c)$$

The matrix  $S_1$  will be so chosen that  $\lambda$  is diagonal, with diagonal matrix elements that will be called  $\lambda_n$ . (This is possible with a constant matrix  $S_1$  since as mentioned above,  $M'$  is taken constant in each range segment). Next, the derivative terms  $\vec{\Gamma}'$  which now multiply the diagonal matrix  $\lambda$ , may be eliminated by standard methods. We make the transformation

$$\vec{\Gamma} = \sigma \vec{f} \quad (22a)$$

where  $\sigma$  is taken to be a diagonal matrix with elements  $\sigma_n$ .

Substituting Eq. (22a) into Eq. (20) and equating the coefficients of  $\gamma_n'$  to zero, we find

$$\sigma_n = e^{-\lambda_n \varphi} ; \quad (22b)$$

constant factors in Eq. (22b) may all be chosen equal to unity.

The transformed range equations are still coupled (although without first derivatives), and have the form

$$d^2 \vec{\gamma} / d\varphi^2 - \lambda^2 \vec{\gamma} + (\tilde{K}^2 + \tilde{W}) \vec{\gamma} = S_1^{-1} \vec{U}. \quad (23)$$

The source term is unaffected by the transformation of Eq. (22a) since  $\vec{U}$  contains a delta function in  $\varphi$ , and  $\sigma(\varphi=0)$  is the identity matrix. The elements of the non-diagonal matrices  $\tilde{K}^2$  and  $\tilde{W}$  are given by

$$\begin{aligned} \tilde{K}_{nm}^2 &= e^{(\lambda_n - \lambda_m)\varphi} K_{nm}^2, \\ \tilde{W}_{nm} &= e^{(\lambda_n - \lambda_m)\varphi} W_{nm}. \end{aligned} \quad (24a)$$

In the spirit of our approximation procedure, these will again be replaced by their values at the midpoints of the  $\varphi$ -intervals, rendering the expression

$$T = \tilde{K}^2 + \tilde{W} - \lambda^2 \quad (25)$$

a constant matrix. The ensuing equation

$$d^2 \vec{\gamma} / d\varphi^2 + T \vec{\gamma} = S_1^{-1} \vec{U} \quad (26)$$

can be decoupled with one final transformation:

$$\vec{\gamma} = S_2 \vec{P} \quad (27)$$

where  $S_2$  is chosen such that

$$S_2^{-1} T S_2 \equiv \Lambda \quad (28)$$

is a diagonal matrix with elements  $\Lambda_n$ . This has finally brought us to the uncoupled system of linear equations with constant coefficients

$$d^2 \vec{P} / d\varphi^2 + \Lambda \vec{P} = \vec{U}^* , \quad (29)$$

where the new source term is

$$\vec{U}^* = (S_1 S_2)^{-1} \vec{U} . \quad (29b)$$

Its general solution is in the source-free region:

$$\vec{P} = \sum_{i=1}^N \vec{\pi}_i (\alpha_i \cos q_i \varphi + \beta_i \sin q_i \varphi), \quad (30a)$$

where  $N$  is the total number of modes,

$$q_i = \Lambda_i^{1/2} , \quad (30b)$$

$\vec{\pi}_i$  are the basis vectors with components

$$\pi_{im} = \delta_{im} , \quad (30c)$$

and  $\alpha_i, \beta_i$  are  $2N$  arbitrary coefficients. If we call the overall transformation matrix

$$U = S_1 \sigma S_2 , \quad (31a)$$

our original range functions of Eq. (15) are obtained as

$$\vec{f} = U \vec{P} , \quad (31b)$$

with components

$$f_n(\rho) = \sum_m U_{nm} (\alpha_m \cos q_m \rho + \beta_m \sin q_m \rho). \quad (31c)$$

The Green's function of Eq. (29a) which describes the behavior of the solution at the source can be obtained from Eq. (31c) in the well-known way,<sup>3-5</sup> see below.

As the last step, we shall have to carry out a matching of the range functions and their derivatives at the range interval boundaries, see Eq. (17). In the last interval which reaches out to infinity, we assume no range dependence so that  $k_p(\rho) = \text{constant} = k_p(\infty)$ , and have the solution<sup>5</sup>

$$f_p^\infty(\rho) = A_p \rho^{1/2} H_0^{(1)}(k_p(\infty) \rho) \quad (32a)$$

(no mode coupling takes place here), comprising only outgoing waves with undetermined modal amplitudes  $A_p$ . In the general (ith) interval, we adopt Eq. (31c) with superscripts  $i$  on  $f_n$ ,  $U_{nm}$ ,  $\alpha_m$ ,  $\beta_m$  and  $q_m$ ; in matrix notation, this gives

$$\vec{f}^i(\rho) = U^i \vec{\mathcal{J}}^i(\rho), \quad (32b)$$

where  $\vec{\mathcal{J}}^i(\rho)$  is a column vector with components

$$\mathcal{J}_q^i(\rho) = \alpha_q^i \cos q_q^i \rho + \beta_q^i \sin q_q^i \rho. \quad (32c)$$



Introducing the diagonal matrix A with diagonal elements  $A_p$ , and the vector  $\vec{H}(\rho)$  with components  $H_p(\rho) = \rho^{1/2} H_0^{(1)}(k_p(\infty)\rho)$ , Eq. (32a) reads in matrix form

$$\vec{f}^\infty(\rho) = A \vec{H}(\rho). \quad (32d)$$

If we designate the last range interval before the infinite interval by  $i=M$ , the matching conditions are (at the boundary  $\rho=\rho_M$ ):

$$\begin{aligned} U^M \vec{f}^M(\rho_M) &= A \vec{H}(\rho_M) \\ U^M \vec{f}^{M'}(\rho_M) &= A \vec{H}'(\rho_M). \end{aligned} \quad (33)$$

We introduce the two-component vectors (2D space)

$$\vec{\alpha}_2^M = \begin{pmatrix} \alpha_2^M \\ \beta_2^M \end{pmatrix}, \quad \vec{H}_p(\rho_M) = \begin{pmatrix} H_p(\rho_M) \\ H_p'(\rho_M) \end{pmatrix} \quad (34a)$$

and the 2 x 2 matrix

$$\mathcal{C}_2^{M\rho_M} = \begin{pmatrix} \cos q_2^M \rho_M & \sin q_2^M \rho_M \\ -q_2^M \sin q_2^M \rho_M & q_2^M \cos q_2^M \rho_M \end{pmatrix} \quad (34b)$$

with inverse

$$(\mathcal{C}_2^{M\rho_M})^{-1} = \begin{pmatrix} \cos q_2^M \rho_M & -\sin q_2^M \rho_M / q_2^M \\ \sin q_2^M \rho_M & \cos q_2^M \rho_M / q_2^M \end{pmatrix}. \quad (34c)$$

Then, Eqs. (33) are solved by

$$\vec{\alpha}_2^M = \sum_p (U_M^{-1} A)_{2p} (\mathcal{C}_2^{M\rho_M})^{-1} \vec{H}_p(\rho_M). \quad (35a)$$

The notation may be further compacted by introducing hypervectors  $\vec{\alpha}^M, \vec{H}(\rho_M)$  whose components (each corresponding to a mode) are the 2-component vectors  $\vec{\alpha}_q^M, \vec{H}_q(\rho_M)$ , respectively, as well as diagonal hypermatrices  $\tilde{\mathcal{E}}$  whose diagonal elements are the  $2 \times 2$  matrices  $\mathcal{E}_q$  (so that the elements of the diagonal hypermatrix  $\tilde{\mathcal{E}}^{-1}$  are the  $2 \times 2$  matrices  $\mathcal{E}_q^{-1}$ ). Eq. (35a) then reads

$$\vec{\alpha}^M = (\tilde{\mathcal{E}}^M \rho_M)^{-1} U_M^{-1} A \vec{H}(\rho_M), \quad (35b)$$

where the previous mode-space matrices  $U_M^{-1}$  and  $A$  are understood as being diagonal in 2D space.

Applying Eq. (17) for the match between intervals  $i$  and  $i+1$ , we have

$$\begin{aligned} U^i \vec{\mathcal{J}}^i(\rho_i) &= U^{i+1} \vec{\mathcal{J}}^{i+1}(\rho_i) \\ U^i \vec{\mathcal{J}}^{i'}(\rho_i) &= U^{i+1} \vec{\mathcal{J}}^{i+1'}(\rho_i) \end{aligned} \quad (36)$$

which, in our hyperspace, is solved by

$$\vec{\alpha}^i = (\tilde{\mathcal{E}}^i \rho_i)^{-1} U_i^{-1} U^{i+1} \tilde{\mathcal{E}}^{i+1, \rho_i} \vec{\alpha}^{i+1}. \quad (37)$$

Finally, we match at the boundary between intervals  $i=1$  (source interval) and  $i=2$ . In the former, there is no coupling, and the solution is<sup>5</sup>

$$f_p(\rho) = A_p^{(1)} \rho^{1/2} [\alpha_p^{(1)} H_0^{(1)}(k_p^{(1)} \rho) + \beta_p^{(1)} H_0^{(2)}(k_p^{(1)} \rho)] \quad (38)$$

where  $A_p^{(0)}$  are  $\frac{un}{\sqrt{k}}$  known amplitude factors which we again arrange in a diagonal matrix  $A^{(0)}$ . Introducing

$$H_p^{(1,2)}(\xi) \equiv \xi^{1/2} H_0^{(1,2)}(k_p^{(0)} \xi) \quad (39a)$$

and forming the known matrix

$$\mathcal{H}_p = \begin{pmatrix} H_p^{(1)}(\rho_1) & H_p^{(2)}(\rho_1) \\ H_p^{(1)'}(\rho_1) & H_p^{(2)'}(\rho_1) \end{pmatrix}, \quad (39b)$$

we obtain from Eqs. (17):

$$A^{(0)} \vec{\alpha}^1 = \tilde{\mathcal{H}}^{-1} U_2 \tilde{\mathcal{E}}^{2, \rho_1} \vec{\alpha}^2, \quad (40)$$

where hypermatrix notation was introduced as before. In this way, all coefficients  $\alpha_q^i$  and  $\beta_q^i$  have been determined; they are in the  $i$ th interval:

$$\vec{\alpha}^i = (\tilde{\mathcal{E}}^{i, \rho_i})^{-1} U_i^{-1} U^{i+1} \tilde{\mathcal{E}}^{i+1, \rho_i} (\tilde{\mathcal{E}}^{i+1, \rho_{i+1}})^{-1} U_{i+1}^{-1} U^{i+2} \tilde{\mathcal{E}}^{i+2, \rho_{i+1}} \dots (\tilde{\mathcal{E}}^{M, \rho_M})^{-1} U_M^{-1} A \vec{H}(\rho_M), \quad (41a)$$

and in the source interval:

$$A^{(0)} \vec{\alpha}^1 = \tilde{\mathcal{H}}^{-1} U_2 \tilde{\mathcal{E}}^{2, \rho_1} (\tilde{\mathcal{E}}^{2, \rho_2})^{-1} U_2^{-1} U^3 \tilde{\mathcal{E}}^{3, \rho_2} \dots (\tilde{\mathcal{E}}^{M-1, \rho_{M-1}})^{-1} U_{M-1}^{-1} U^M \tilde{\mathcal{E}}^{M, \rho_{M-1}} (\tilde{\mathcal{E}}^{M, \rho_M})^{-1} U_M^{-1} A \vec{H}(\rho_M). \quad (41b)$$

These expressions still contain the unknown coefficients  $A^{(0)}$  and  $A$ ,

but as was shown earlier,<sup>5</sup> matching to the source leads to

$$A_p^{(0)} = \frac{\mu_p(z_0, 0)}{4i(\alpha_p^2 - \beta_p^2)} \frac{\varphi(z_0)}{\rho_w}, \quad (42)$$

so that in principle, the remaining unknowns  $A_p$  (the outgoing mode amplitudes at infinity) can be determined by the requirement that Eqs. (41b) and (42) be satisfied. The implementation of this prescription, however, requires an iterative approach which will be described below.

### III. Numerical Results and Discussion

The above procedure will now be applied to the calculation of transmission loss in some examples of channels with simple and also with realistic velocity profiles. In order to be able to handle the latter case, we shall use the previously employed method<sup>5</sup> for calculating local depth functions  $u_n(z, \varphi)$  and their eigenvalues  $k_n^2(\varphi)$  in a numerically given profile, which consists in dividing the ocean into (horizontally bounded) depth segments or layers, in each of which  $k^2(z, \varphi)$  is approximated linearly, so that the depth functions are given by Airy functions. The layer below the (locally flat) ocean floor is simply modeled as a liquid with constant bottom density  $\rho_B$ , and with a  $z$ -independent local sound velocity  $c_B(\varphi)$ , which reaches down to  $z \rightarrow \infty$ . The eigenvalues  $k_n^2(\varphi)$  are then obtained by satisfying the boundary conditions of a pressure release surface at the ocean surface  $z=z_s$ , i.e.,



$u_p(0, \vec{\xi}) = 0$ , and of an exponentially decaying solution in the bottom layer. Further details are given in Reference 5.

REFERENCES

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The following computer program is based on the foregoing theory; it computes range functions and solves the coupled-mode program of sound propagation in a range dependent ocean.

The main program GPC (pp. 24 - 32) computes the propagation loss as a function of range and depth. It calls all the subroutines; it calculates the range functions, it calculates the coefficients of the linear combinations of the range functions in all the range segments.

The subroutine NAGL (p.33) interfaces the range function and depth function programs; it calls the depth function program which is described in Volume I of this report.

The subroutine FORM (pp.34-37) calculates the U's and the q-eigenvalues

The subroutine COEF1 (pp.38-40) calculates the mode coefficients in the source segment and in the infinite segment by matching directly (since all the internal boundary conditions are here lumped together into one matrix)

The subroutine HMDIAG (p.41) diagonalizes an antihermitean matrix (by relating it to the diagonalization of a Hermitean one)

The subroutine MXEL (p.42) computes integrals  $\int u_n u_m z \, dz$  containing Airy functions

The subroutine HANK (p.43) calls the Hankel function subroutine

The subroutine HANGL (p.43) computes the Hankel functions of argument X

The subroutine DAIRY (p.44-50) is the same as in the depth function program (Vol. I of this report)

The subroutine PLT (p.51-52) is a plot routine to plot the propagation loss

The last page (p.53) tells what has to be loaded into the range, depth, and plotting programs.



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TY GPC

```

COMMON /DEPTH/ XKN(15,4),
1UO(15)
COMMON /FORMC/ X(1,2,4),ZB(4),NM,NL,NLM,NLP,XK(4,4),DEN(4)
1,XNAG(4)
  DIMENSION XKT(4),X(15,4),
1U(15,15),UI(15,15),R(15),AM(30,30),V(15),DUM(30,30),
2A(15),AO(15),ALF(3,15),BET(3,15),US(2,15,15),UIS(2,15,15),
3QS(2,15),FR(15),UR(15),IDUT(15),G(30),XKNT(15),NMODE(4)
  COMPLEX U,UI,Q,AM,SUM,EYE,V,DUM,A,AO,ALF,BET,US,UIS,QS,SUM2,
1FR,UN,C1,S1,H1,H1P,H2,H2P,A1,A2,A3,A4,A5,TEM,QR
  DIMENSION UOD(15),URD(15,4),XKND(15,4)
  DOUBLE PRECISION XNAG,UOD,URD,XKND
  EYE=(0.,1.)
  PI=3.14159
  TYPE 720
720 FORMAT(' TYPE 0 FOR NO COUPLING, 1 FOR COUPLING')
  ACCEPT 2,ICOUP
  TYPE 1
  1 FORMAT(' INPUT NUMBER OF MODES, LAYERS, AND SEGMENTS')
  ACCEPT 2,NM,NL,NS
  TYPE 2,NM,NL,NS
  2 FORMAT(5G)
  NLM=NL-1
  NLP=NL+1
  NSM=NS-1
  NSP=NS+1
  TYPE 460
460 FORMAT(' INPUT NUMBER OF MODES FOR EACH PROFILE')
  ACCEPT 2,(NMODE(I),I=1,NSP)
  TYPE 2,(NMODE(I),I=1,NSP)
  TYPE 6
  6 FORMAT(' INPUT FREQUENCY')
  ACCEPT 2,F
  TYPE 2,F
  F=2.*PI*F
  TYPE 700
700 FORMAT(' INPUT CONVERGENCE CRITERION')
  ACCEPT 2,FRAC
  TYPE 2,FRAC
  TYPE 3
  3 FORMAT(' INPUT SEGMENT POSITIONS')
  X(1)=2000.
  X(2)=20000.
  X(3)=40000.
  TYPE 2,(X(I),I=1,NS)
  ACCEPT 2,(X(I),I=1,NS)
  XNAG(1)=0.
  DO 200 I=1,NS
200 XNAG(I+1)=X(I)
  TYPE 4
  4 FORMAT(' INPUT SOURCE DEPTH, RECEIVER DEPTH, NL-1 PROFILE',
1' DEPTHS, NS+1 BOTTOM DEPTHS')
  ACCEPT 2,ZS,ZR,(Z(I),I=2,NL),(ZB(I),I=1,NSP)
  TYPE 2,ZS,ZR,(Z(I),I=2,NL),(ZB(I),I=1,NSP)
  Z(1)=0.
  DO 101 IZS=1,NL
  IF(ZS.LT.Z(IZS+1)) GO TO 9
101 CONTINUE
  8 TYPE 5
  5 FORMAT(' INPUT SOUND SPEED PROFILE FOR EACH SEGMENT')
  DO 802 IS=1,NSP
802 XK(1,IS)=1520.
  XK(2,1)=1515.
  XK(2,2)=1514.5
  XK(2,3)=1510.

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XK(2,4)=1505.
XK(3,1)=1490.
XK(3,2)=1490.
XK(3,3)=1490.
XK(3,4)=1490.
XK(4,1)=1540.
XK(4,2)=1540.
XK(4,3)=1540.
XK(4,4)=1540.
DO 100 IS=1,NSP
C   ACCEPT 2,(XK(IL,IS),IL=1,NLP)
   TYPE 2,(XK(IL,IS),IL=1,NLP)
DO 100 IL=1,NLP
XK(IL,IS)=(F/XK(IL,IS))*2
100 CONTINUE
   TYPE 7
   7 FORMAT(' INPUT RELATIVE DENSITIES FOR EACH LAYER')
DO 803 IL=1,NLP
803 DEN(IL)=1.
   TYPE 2,(DEN(IL),IL=1,NLP)
C   ACCEPT 2,(DEN(IL),IL=1,NLP)
   TYPE 40
40 FORMAT(' TYPE NUMBER OF MODES TO BE OUTPUT')
   ACCEPT 2,NOUT
   TYPE 2,NOUT
   NOUT2=2*NOUT
   IF(NOUT.EQ.0) GO TO 41
   IF(NOUT.EQ.NM) GO TO 43
   TYPE 42
42 FORMAT(' TYPE WHICH MODES ARE TO BE OUTPUT')
   ACCEPT 2,(IOUT(I),I=1,NOUT)
   GO TO 41
43 DO 44 I=1,NM
44 IOUT(I)=I
C
C   DETERMINE DEPTH FUNCTIONS
C
41 CALL NAGL(NM,NLP,NSP,NMODE,Z,ZB,XK,XKN,UO,UR,ZS,XKND,UOD,URD)
C
C   FORM PRODUCT OF MATRICES ON PAGE 44...U2*H2(X1)*H2I(X2)*U2I*U3*...
C   FOR USE IN CONNECTING SOURCE INTERVAL WITH INFINITE INTERVAL.
C   SIZE OF MATRIX IS NOW TWICE THE NUMBER OF MODES TO ALLOW FOR
C   INGOING AND OUTGOING SOLUTIONS. MATRIX U IS DIAGONAL IN THE 2
C   SOLUTIONS. MATRIX SCRIPT-C IS DIAGONAL IN MODE NUMBER.
C
NM2=2*NM
DO 106 I=2,NM2
JM=I-1
DO 107 J=1,JM
AM(I,J)=0.
107 AM(J,I)=0.
106 AM(I,I)=1.
AM(1,1)=1.
C   TYPE 201,((AM(I,J),J=1,NM2),I=1,NM2)
201 FORMAT(' 201 AM'/(12G))
DO 108 IS=2,NS
ISM=IS-1
ISTEM=IS
CALL FORM(ISTEM,U,UI,Q,ICOU,NM2)
C   TYPE 621,(G(I),I=1,NM)
621 FORMAT(' 621 Q'/(6G))
DO 113 I=1,NM
DO 110 J=1,NM
US(ISM,I,J)=U(I,J)
113 US(ISC(I),I,J)=UI(I,J)
DO 119 I=1,NM

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119 QS(ISM,I)=Q(I)
C
C      MULTIPLY FROM RIGHT BY U
C
      DO 109 I=1,NM2
      DO 109 K=1,NM
      DO 109 J=1,2
      SUM=0.
      DO 110 L=1,NM
110 SUM=SUM+AM(I,2*(L-1)+J)*U(L,K)
109 DUM(I,2*(K-1)+J)=SUM
C      TYPE 222,((DUM(I,J),J=1,NM2),I=1,NM2)
222 FORMAT(' 222 DUM'/(12G))
C
C      MULTIPLY FROM RIGHT BY SCRIPT-C EVALUATED AT LEFT OF SEGMENT
C
      DO 111 L=1,NM2
      DO 111 I=1,NM
      I1=2*(I-1)
      C1=CCOS(Q(I)*X(ISM))
      S1=CSIN(Q(I)*X(ISM))
      AM(L,I1+1)=DUM(L,I1+1)*C1-DUM(L,I1+2)*Q(I)*S1
111 AM(L,I1+2)=DUM(L,I1+1)*S1+DUM(L,I1+2)*Q(I)*C1
C      TYPE 223,((AM(I,J),J=1,NM2),I=1,NM2)
223 FORMAT(' 223 AM'/(12G))
C
C      MULTIPLY FROM RIGHT BY SCRIPT-C INVERSE EVALUATED AT RIGHT OF
C      SEGMENT
C
      DO 112 L=1,NM2
      DO 112 I=1,NM
      I1=2*(I-1)
      C1=CCOS(Q(I)*X(IS))
      S1=CSIN(Q(I)*X(IS))
      DUM(L,I1+1)=AM(L,I1+1)*C1+AM(L,I1+2)*S1
112 DUM(L,I1+2)=(-AM(L,I1+1)*S1+AM(L,I1+2)*C1)/Q(I)
C      TYPE 224,((DUM(I,J),J=1,NM2),I=1,NM2)
224 FORMAT(' 224 DUM'/(12G))
C
C      MULTIPLY FROM RIGHT BY U INVERSE
C
      DO 113 I=1,NM2
      DO 113 K=1,NM
      DO 113 J=1,2
      SUM=0.
      DO 114 L=1,NM
114 SUM=SUM+DUM(I,2*(L-1)+J)*UI(L,K)
113 AM(I,2*(K-1)+J)=SUM
C      TYPE 225,((AM(I,J),J=1,NM2),I=1,NM2)
225 FORMAT(' 225 AM'/(12G))
109 CONTINUE
C
C      MULTIPLY FROM LEFT BY SCRIPT-H INVERSE
C
      X2=SQRT(X(1))
      DO 115 I=1,NM
      I1=2*(I-1)
      RHO=SQRT(XKN(I,1))*X(1)
      CALL HANK(RHO,H1,H1P,H2,H2P)
C      TYPE 243,I,RHO,H1,H1P,H2,H2P
      XNTEM=RHO/X(1)
      A1=X2*H1
      A2=X2*H2
      A3=XNTEM*X2*H1P*H1/(2.*X2)
      A4=XNTEM*X2*H2P*H2/(2.*X2)
      A5=A1+A4-A2+A3

```

```

TEN=A1
A1=A4/A5
A4=TEM/A5
A2=-A2/A5
A3=-A3/A5
DO 115 L=1,NM2
DUM(I1+1,L)=A1*AM(I1+1,L)+A2*AM(I1+2,L)
115 DUM(I1+2,L)=A3*AM(I1+1,L)+A4*AM(I1+2,L)
C   TYPE 226,((DUM(I,J),J=1,NM2),I=1,NM2)
226 FORMAT(' 226 DUM'/(12G))
C
C   MULTIPLY FROM RIGHT BY H
C
XNS=SQRT(X(NS))
DO 116 J=1,NM
J1=2*(J-1)
RHO=SQRT(XKN(J,NSP))*X(NS)
CALL HANK(RHO,H1,H1P,H2,H2P)
C   TYPE 243,J,RHO,H1,H1P,H2,H2P
XKTEM=RHO/X(NS)
DO 116 I=1,NM2
116 AM(I,J)=DUM(I,J1+1)*XNS*H1+DUM(I,J1+2)*(XKTEM*XNS*H1P+H1/(2.*
1XNS))
C   TYPE 227,((AM(I,J),J=1,NM),I=1,NM2)
227 FORMAT(' 227 AM'/(6G))
C
C   FORM MATRIX TO BE USED IN DETERMINING COEFFICIENTS IN SOURCE
C   SEGMENT
C
DO 117 I=1,NM
117 V(I)=UO(I)*DEN(IZS)/(4.*EYE)
C   TYPE 228,(V(I),I=1,NM)
228 FORMAT(' 228 V'/(6G))
C
C   FIND COEFFICIENTS FOR SOURCE INTERVAL
C
CALL COEF1(NM,AM,V,A,A0,ALF,BET,FRAC)
C
C   DETERMINE COEFFICIENTS FOR OTHER INTERVALS USING RELATION ON
C   PAGE 44
C
DO 120 I=2,NM2
JM=I-1
DO 121 J=1,JM
AM(I,J)=0.
121 AM(J,I)=0.
120 AM(I,I)=1.
AM(1,1)=1.
C   TYPE 240,((AM(I,J),J=1,NM2),I=1,NM2)
240 FORMAT(' 240 AM'/(12G))
C
C   MULTIPLY FROM RIGHT BY SCRIPT-C INVERSE EVALUATED AT BOUNDARY OF
C   INFINITE SEGMENT
C
DO 122 L=1,NM2
DO 122 I=1,NM
I1=2*(I-1)
C1=CCOS(QS(NSH,I))*X(NS)
S1=CSIN(QS(NSH,I))*X(NS)
DUM(L,I1+1)=AM(L,I1+1)*C1+AM(L,I1+2)*S1
122 DUM(L,I1+2)=(-AM(L,I1+1)*S1+AM(L,I1+2)*C1)*QS(NSH,I)
C   TYPE 241,((DUM(I,J),J=1,NM2),I=1,NM2)
241 FORMAT(' 241 DUM'/(12G))
C
C   MULTIPLY FROM RIGHT BY U INVERSE
C

```



```

DO 123 I=1,NM2
DO 123 K=1,NM
DO 123 J=1,2
SUM=0.
DO 124 L=1,NM
124 SUM=SUM+DUM(I,2*(L-1)+J)*UIS(NSM,L,I)
123 AM(I,2*(K-1)+J)=SUM
C   TYPE 242,((AM(I,J),J=1,NM2),I=1,NM2)
242 FORMAT(' 242 AM'/(12G))
C
C   MULTIPLY FROM RIGHT BY H EVALUATED AT BOUNDARY OF INFINITE
C   SEGMENT
C
XNS=SQRT(X(NS))
DO 125 J=1,NM
J1=2*(J-1)
RHO=SQRT(XKN(J,NSP))*X(NS)
CALL HANK(RHO,H1,H1P,H2,H2P)
C   TYPE 243,J,RHO,H1,H1P,H2,H2P
243 FORMAT(' 243 J,RHO,H1,H1P,H2,H2P'/(10G))
XKTEM=RHO/X(NS)
DO 125 I=1,NM2
125 DUM(I,J)=AM(I,J1+1)*XNS*H1+AM(I,J1+2)*(XKTEM*XNS*H1P+H1/(2.*XNS))
C   TYPE 244,((DUM(I,J),J=1,NM),I=1,NM2)
244 FORMAT(' 244 DUM'/(6G))
C
C   MULTIPLY FROM RIGHT BY COEFFICIENTS FOR INFINITE SEGMENT TO DEFINE
C   COEFFICIENTS FOR LAST FINITE SEGMENT
C
DO 126 I=1,NM
I1=2*(I-1)
SUM=0.
SUM2=0.
DO 127 J=1,NM
SUM=SUM+DUM(I1+1,J)*A(J)
127 SUM2=SUM2+DUM(I1+2,J)*A(J)
AM(I1+1,1)=SUM
AM(I1+2,1)=SUM2
ALF(NS,I)=SUM
126 BET(NS,I)=SUM2
C   TYPE 245,(AM(I,1),I=1,NM2)
245 FORMAT(' 245 AM'/(12G))
C   TYPE 246,(ALF(NS,I),I=1,NM)
246 FORMAT(' 246 ALF'/(6G))
C   TYPE 247,(BET(NS,I),I=1,NM)
247 FORMAT(' 247 BET'/(6G))
C
C   FIND COEFFICIENTS OF OTHER SEGMENTS BY MATCHING SEQUENTIALLY...
C   USE RELATION IN MIDDLE OF PAGE 43
C
NS2=NS-2
IF(NS2.LT.1) GO TO 20
IS=NS
DO 128 IS2=1,NS2
IS=IS-1
ISM=IS-1
C
C   MULTIPLY FROM LEFT BY SCRIPT S FOR PREVIOUS SEGMENT EVALUATED
C   AT BOUNDARY BETWEEN CURRENT AND PREVIOUS SEGMENTS
C
DO 129 I=1,NM
I1=2*(I-1)
C1=CCCC(QS(IS,I))*X(IS)
C1=CCCC(QS(IS,I))*X(ISO)
DUM(I1+1,1)=C1*AM(I1+1,1)+S1*AM(I1+2,1)
129 DUM(I1+2,1)=QS(IS,I)*(-S1*AM(I1+1,1)+C1*AM(I1+2,1))

```

```

C      TYPE 248,(DUM(I,1),I=1,NM2)
248 FORMAT(' 248 DUM'/12G)
C
C      MULTIPLY FROM LEFT BY U FOR PREVIOUS SEGMENT
C
      DO 130 I=1,NM
      I1=2*(I-1)
      DO 130 J=1,2
      SUM=0.
      DO 131 K=1,NM
131 SUM=SUM+US(IS,I,K)*DUM(2*(K-1)+J,1)
130 AM(I1+J,1)=SUM
C      TYPE 249,(AM(I,1),I=1,NM2)
249 FORMAT(' 249 AM'/12G)
C
C      MULTIPLY FROM LEFT BY U INVERSE FOR CURRENT SEGMENT
C
      DO 132 I=1,NM
      I1=2*(I-1)
      DO 132 J=1,2
      SUM=0.
      DO 133 K=1,NM
133 SUM=SUM+UIS(ISM,I,K)*AM(2*(K-1)+J,1)
132 DUM(I1+J,1)=SUM
C      TYPE 250,(DUM(I,1),I=1,NM2)
250 FORMAT(' 250 DUM'/12G)
C
C      OBTAIN COEFFICIENTS FOR CURRENT SEGMENT BY MULTIPLYING FROM
C      LEFT BY SCRIPT-C INVERSE FOR CURRENT SEGMENT EVALUATED AT
C      RIGHT OF SEGMENT
C
      DO 134 I=1,NM
      I1=2*(I-1)
      C1=CCOS(QS(ISM,I)*X(IS))
      S1=CSIN(QS(ISM,I)*X(IS))
      SUM=C1*DUM(I1+1,1)-S1*DUM(I1+2,1)/QS(ISM,I)
      SUM2=S1*DUM(I1+1,1)+C1*DUM(I1+2,1)/QS(ISM,I)
      AM(I1+1,1)=SUM
      AM(I1+2,1)=SUM2
      ALF(IS,I)=SUM
134 BET(IS,I)=SUM2
C      TYPE 251,(AM(I,1),I=1,NM2)
251 FORMAT(' 251 AM'/12G)
C      TYPE 252,(ALF(IS,I),I=1,NM)
252 FORMAT(' 252 ALF'/6G)
C      TYPE 253,(BET(IS,I),I=1,NM)
253 FORMAT(' 253 BET'/6G)
128 CONTINUE
C
C      RECOMPUTE COEFFICIENTS IN SOURCE SEGMENT.
C
C      TYPE 620,(QS(1,I),I=1,NM)
620 FORMAT(' 620 QS'/6G)
      DO 600 I=1,NM
      I1=2*(I-1)
      C1=CCOS(QS(1,I)*X(1))
      S1=CSIN(QS(1,I)*X(1))
      DUM(I1+1,1)=C1*AM(I1+1,1)+S1*AM(I1+2,1)
600 DUM(I1+2,1)=QS(1,I)*(-S1*AM(I1+1,1)+C1*AM(I1+2,1))
C      TYPE 601,(DUM(I,1),I=1,NM2)
601 FORMAT(' 601 DUM'/12G)
      DO 602 I=1,NM
      I1=2*(I-1)
      DO 602 J=1,2
      SUM=0.
      DO 603 K=1,NM

```

```

603 SUM=SUM+US(1,I,N)*DUM(2*(K-1)+J,1)
602 AM(I1+J,1)=SUM
C   TYPE 604,(AM(I,1),I=1,NM2)
604 FORMAT(' 604 AM'/12G)
X2=SQRT(X(1))
DO 605 I=1,NM
I1=2*(I-1)
RHO=SQRT(XNM(I,1))*X(1)
CALL HANK(RHO,H1,H1P,H2,H2P)
XNTEM=RHO/X(1)
A1=X2*H1
A2=X2*H2
A3=XNTEM*X2*H1P*H1/(2.*X2)
A4=XNTEM*X2*H2P*H2/(2.*X2)
A5=A1*A4-A2*A3
TEM=A1
A1=A4/A5
A4=TEM/A5
A2=-A2/A5
A3=-A3/A5
DUM(I1+1,1)=(A1*AM(I1+1,1)+A2*AM(I1+2,1))/A0(I)
605 DUM(I1+2,1)=(A3*AM(I1+1,1)+A4*AM(I1+2,1))/A0(I)
C   TYPE 606,(DUM(I,1),I=1,NM2)
606 FORMAT(' 606 DUM'/12G)
C
C   INPUT RANGES AND DEPTHS FOR FINDING PROP LOSS
C
20 TYPE 21
21 FORMAT(' INPUT 0 FOR RANGE VARIATION, 1 FOR DEPTH VARIATION')
ACCEPT 2,IRZ
IF(IRZ.EQ.1) GO TO 22
TYPE 23
23 FORMAT(' INITIAL RANGE, FINAL RANGE, '
1' STEP SIZE')
ACCEPT 2,R1,R2,DR
R1=R1*1000.
R2=R2*1000.
DR=DR*1000.
R1=R1-DR
R=R1
NR=(R2-R)/DR+.01
NZ=1
DZ=0.
GO TO 24
22 TYPE 25
25 FORMAT(' INPUT RANGE, INITIAL DEPTH, FINAL DEPTH, STEP SIZE')
ACCEPT 2,R,Z1,Z2,DZ
Z1=Z1-DZ
NZ=(Z2-Z1)/DZ+.01
NR=1
DR=0.
24 DO 135 IR=1,NR
R=R+DR
ROUT=R/1000.
DO 136 IS=1,NS
IF(R.LT.X(IS)) GO TO 925
136 CONTINUE
C
C   RANGE IS IN INFINITE SEGMENT
C
IS=NSP
ZB1=SQRT(ZB(1))
ZBN=SQRT(ZB(NSP))
DO 137 I=1,NM
RHO=SQRT(XNM(I,NSP))*R
CALL HANK(RHO,H1,H1P,H2,H2P)

```

```

137 FR(I)=A(I)*H1/ZB1
    SUM=0.
    DO 750 I=1,NM
750 SUM=SUM+FR(I)*UR(I,NCP)
    SUM=SUM/ZBN
    GO TO 24
825 IF(IS.EQ.1) GO TO 27
C
C
C    RANGE IS NOT IN EITHER SOURCE OR INFINITE SEGMENT
    ISM=IS-1
    X2=SQRT(R)
    ZB1=SQRT(ZB(1))
    DO 139 I=1,NM
    SUM=0.
    DO 139 J=1,NM
    QR=QS(ISM,J)*R
139 SUM=SUM+US(ISM,I,J)*(ALF(IS,J)*CCOS(QR)+BET(IS,J)*CSIN(QR))/X2
138 FR(I)=SUM/ZB1
    GO TO 26
C
C
C    RANGE IS IN SOURCE SEGMENT
27 ZB1=SQRT(ZB(1))
    DO 140 I=1,NM
    RHO=SQRT(XKN(I,1))*R
    CALL HANK(RHO,H1,H1F,H2,H2F)
140 FR(I)=A0(I)*(ALF(1,I)*H1+BET(1,I)*H2)/ZB1
C
C
C    OUTPUT MODE INFORMATION
26 IF(NOUT.EQ.0) GO TO 45
    J=1
    DO 150 I=1,NM
    IF(I.NE.IDOUT(J)) GO TO 150
    I1=2*(J-1)
    G(I1+1)=CABS(FR(I))
    G(I1+2)=CLOG(FR(I)/G(I1+1))/EYE*180./PI
    IF(J.EQ.NOUT) GO TO 46
    J=J+1
150 CONTINUE
46 WRITE (23,47) ROUT,(G(I),I=1,NOUT2)
47 FORMAT(F10.3,1P6E10.3/(10X,1P6E10.3))
C
C
C    COMPUTE PROP LOSS
45 IF(IS.EQ.NSP) GO TO 751
    ISP=IS+1
    FRAC=(R-XNAG(IS))/(XNAG(ISF)-XNAG(IS))
    SUM=0.
    ZB1=SQRT(ZB(IS))
    ZB2=SQRT(ZB(ISF))
    DO 142 I=1,NM
142 SUM=SUM+FR(I)*(UR(I,IS)/ZB1+FRAC*(UR(I,ISF)/ZB2-UR(I,IS)/ZB1))
751 TL=-20.*ALOG10(4.*PI*CABS(SUM))
    IF(IRZ.EQ.1) GO TO 28
    IF(ROUT.NE.1..OR.ROUT.NE.10..OR.ROUT.NE.25..OR.ROUT.NE.50..OR.ROUT
    1.NE.75..OR.ROUT.NE.100.) GO TO 28
    TYPE 29,ROUT,TL
29 FORMAT(F13.5,F10.2)
1234 CONTINUE
    WRITE (22,999) ROUT,TL
999 FORMAT(F13.5,F10.2)
    GO TO 135
28 TYPE 29,ROUT,TL
    WRITE (22,29) ZD,TL

```



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```
135 CONTINUE  
T=-1.  
WRITE (22,29) T  
TYPE 30  
30 FORMAT(' INPUT 0 FOR NEW RANGE OR DEPTH, 1 FOR END')  
ACCEPT 2,IEND  
IF(IEND.EQ.0) GO TO 20  
STOP  
END
```

TY HACL

SUBROUTINE HACL(M,RLP,NSP,NHGT,Z,ZB,XKNS,IOS,URS,ZG,XKN,UN)

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```

C
TY NAGL
SUBROUTINE NAGL(NM,NLP,NSP,NMODE,Z,ZB,XK,XKNS,UOS,URS,ZS,XKN,UO,
1UR)
DIMENSION Z(4),ZB(4),XK(4,4),XKN(15,4),UO(15),UR(15,4),NMODE(4)
DIMENSION XKNS(15,4),UOS(15),URS(15,4)
DOUBLE PRECISION UO,UR,XKN
DO 100 ISP=1,NSP
IM=NMODE(ISP)
IF(IM.EQ.0) GO TO 300
IF(ISP.EQ.1)READ(25,1)(UO(I),I=1,IM)
1 FORMAT(100D)
C IF(ISP.EQ.1)TYPE 237,(UO(I),I=1,IM)
237 FORMAT(' 237 UO'/(3G))
READ(27,1)(UR(I,ISP),I=1,IM)
C TYPE 238,(UR(I,ISP),I=1,IM)
238 FORMAT(' 238 UR' /3G)
READ(28,1)(XKN(I,ISP),I=1,IM)
DO 101 I=1,IM
101 XKN(I,ISP)=XKN(I,ISP)**2
C TYPE 239,(XKN(I,ISP),I=1,IM)
239 FORMAT(' 239 XKN' /3G)
IF(ZS.GT.0..OR.ISP.NE.1) GO TO 100
DO 302 I=1,IM
302 UO(I)=UR(I,1)
GO TO 100
300 IM=NMODE(1)
NMODE(ISP)=NMODE(1)
DO 301 I=1,IM
UR(I,ISP)=UR(I,1)
301 XKN(I,ISP)=XKN(I,1)
100 CONTINUE
DO 110 ISP=1,NSP
IM=NMODE(ISP)
DO 110 I=1,IM
UOS(I)=UO(I)
URS(I,ISP)=UR(I,ISP)
110 XKNS(I,ISP)=XKN(I,ISP)
RETURN
END

```

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TY FRAXC

.TTY FORM

.TY FORM

```
SUBROUTINE FORM(IS,U,UI,Q,ICOU,MODE)
COMMON /DEPTH/ XKN(15,4),
1UO(15)
COMMON /FORMC/ X(3),Z(4),ZB(4),NM,NL,NLM,NLP,XK(4,4),DEN(4)
1,XNAG(4)
DIMENSION AC(3),ST(15),CT(15),C(15,15),V(15,15),DUM(15,15),
1S(15,15),SI(15,15),EIG(15),U(15,15),UI(15,15),Q(15),DMAT(15,15),
2ENAT(120),DVEC(15,15),DEIG(15),WK(45),WK2(1480),SIG(15,15),
3XMXEL(15,15),MODE(4)
COMPLEX SUN,C,V,DUM,S,SI,EIG,U,UI,Q,DMAT,ENAT,DVEC,SIG
DOUBLE PRECISION XNAG
```

```
C
C COMPUTES U=S*SIGMA*S1 AND ITS INVERSE AND THE EIGENVALUES Q OF
C FINAL TRANSFORMATION
MATSIZ=15
NP=NM*(NM+1)/2
NQ=3*NM
NWK2=2*NM*(NM+1)
ISM=IS-1
ISP=IS+1
```

```
C
C LINEARIZE KN**2 IN RANGE
```

```
C
DO 100 I=1,NM
ST(I)=(XKN(I,ISP)-XKN(I,ISM))/(X(IS)-X(ISM))
100 CT(I)=XKN(I,IS)-ST(I)*X(ISM)
C TYPE 202,(ST(I),I=1,NM)
202 FORMAT(' 202 ST'/3G)
C TYPE 203,(CT(I),I=1,NM)
203 FORMAT(' 203 CT'/3G)
```

```
C
C LINEARIZE K**2 IN DEPTH AND RANGE
```

```
C
DO 101 IL=1,NLM
101 AC(IL)=(XK(IL+1,ISP)-XK(IL,ISP)-(XK(IL+1,IS)-XK(IL,IS)))/
1((X(IS)-X(ISM))*(Z(IL+1)-Z(IL)))
AC(NL)=((XK(NLP,ISP)-XK(NL,ISP))/(ZB(ISP)-Z(NL))
1-(XK(NLP,IS)-XK(NL,IS))/(ZB(IS)-Z(NL)))/(X(IS)-X(ISM))
C TYPE 204,(AC(I),I=1,NL)
204 FORMAT(' 204 AC'/3G)
```

```
C
C FORM MATRIX C ON PAGE 37
```

```
C
CALL MXEL(MODE,IS,XNAG,XMXEL,ZD)
C TYPE 205,((XMXEL(I,J),J=1,NM),I=1,NM)
205 FORMAT(' 205 XMXEL'/3G)
DO 104 I=2,NM
IM=I-1
DO 102 J=1,IM
C(I,J)=-XMXEL(I,J)/((ST(I)-ST(J))*(X(IS)+X(ISM))/2+CT(I)-CT(J))
102 C(J,I)=-C(I,J)
104 C(I,I)=0.
C(I,1)=0.
C TYPE 206,((C(I,J),J=1,NM),I=1,NM)
206 FORMAT(' 206 C'/16G)
```

```
C
C FORM MATRIX V ON PAGE 36
```

```
C
DO 103 I=1,NM
DO 105 J=1,I
SUM=0.
```

```

DO 107 K=1,NM
IF(K.EQ.I.OR.K.EQ.J) GO TO 107
SUM=SUM+C(I,K)*C(K,J)
107 CONTINUE
SUM=-SUM
IF(I.EQ.J) GO TO 2
V(I,J)=-C(I,J)*(ST(I)-ST(J))/((ST(I)-ST(J))*(X(IS)+X(ISM))/2.
1+CT(I)-CT(J))+SUM
V(J,I)=-C(J,I)*(ST(J)-ST(I))/((ST(J)-ST(I))*(X(IS)+X(ISM))/2.
1+CT(J)-CT(I))+SUM
GO TO 105
2 V(I,I)=SUM
105 CONTINUE
IF(ICOUF.EQ.1) GO TO 802
DO 803 I=1,NM
DO 804 J=1,NM
U(I,J)=0.
U(J,I)=0.
UI(I,J)=0.
804 UI(J,I)=0.
U(I,I)=1.
UI(I,I)=1.
Q(I)=SQRT((XKN(I,IS)+XKN(I,ISP))/2.)
803 CONTINUE
C TYPE 805,((U(I,J),J=1,NM),I=1,NM)
805 FORMAT(' 805 U'/(6G))
C TYPE 806,((UI(I,J),J=1,NM),I=1,NM)
806 FORMAT(' 806 UI'/(6G))
C TYPE 807,(Q(I),I=1,NM)
807 FORMAT(' 807 Q'/(6G))
RETURN
802 CONTINUE
C TYPE 207,((V(I,J),J=1,NM),I=1,NM)
207 FORMAT(' 207 V'/(6G))
C
C DIAGONALIZE C.
C C IS ANTISYMMETRIC (ANTIHERMITIAN)...EIGENVALUES ARE IMAGINARY
C
CALL HMDIAG(NM,NP,NQ,C,EIG,S,SI,DMAT,EMAT,DVEC,DEIG,WK)
C TYPE 208,((S(I,J),J=1,NM),I=1,NM)
208 FORMAT(' 208 S'/(6G))
C TYPE 209,((SI(I,J),J=1,NM),I=1,NM)
209 FORMAT(' 209 SI'/(6G))
C TYPE 210,(EIG(I),I=1,NM)
210 FORMAT(' 210 EIG'/(6G))
C
C DEFINE DIAGONAL MATRIX SIGMA
C
DO 108 I=1,NM
SIG(I,I)=CEXP(-EIG(I)*(X(IS)+X(ISM))/2.)
IF(I.EQ.1) GO TO 108
IM=I-1
DO 300 J=1,IM
SIG(I,J)=0.
300 SIG(J,I)=0.
108 CONTINUE
C TYPE 211,((SIG(I,J),J=1,NM),I=1,NM)
211 FORMAT(' 211 SIG'/(6G))
C
C FORM MATRIX V+K
C
DO 109 I=1,NM
109 V(I,I)=V(I,I)*(XKN(I,IS)+XKN(I,ISP))/2.
C TYPE 212,((V(I,J),J=1,NM),I=1,NM)
212 FORMAT(' 212 V'/(6G))
C

```



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```

C
C
C      FORM MATRIX (V+K)*S
      DO 110 I=1,NM
      DO 110 J=1,NM
      SUM=0.
      DO 111 K=1,NM
111  SUM=SUM+V(I,K)*S(K,J)
110  DUM(I,J)=SUM
C      TYPE 213,((DUM(I,J),J=1,NM),I=1,NM)
213  FORMAT (' 213 DUM'/(6G))
C
C      FORM MATRIX T (CALL IT V TO SAVE SPACE) ON PAGE 38
C
      DO 114 I=1,NM
      DO 112 J=1,NM
      SUM=0.
      DO 113 K=1,NM
113  SUM=SUM+SI(I,K)*DUM(K,J)
112  V(I,J)=SUM*CEXP((EIG(I)-EIG(J))*(X(IS)+X(ISM))/2.)
114  V(I,I)=V(I,I)-EIG(I)**2
C      TYPE 214,((V(I,J),J=1,NM),I=1,NM)
214  FORMAT (' 214 V'/(6G))
C
C      DIAGONALIZE T(CALLED V IN CODE)
C
      CALL EIGCC(V,NM,MATSIZ,2,Q,U,MATSIZ,WK2,IER)
C      TYPE 888,(Q(I),I=1,NM)
888  FORMAT (' 888 Q'/(6G))
      DO 450 I=1,NM
450  Q(I)=CSQRT(Q(I))
C      TYPE 215,(Q(I),I=1,NM)
215  FORMAT (' 215 Q'/(6G))
C      TYPE 216,((U(I,J),J=1,NM),I=1,NM)
216  FORMAT (' 216 U'/(6G))
C
C      THE FOLLOWING CARD IS TEMPORARY TO CHECK PERFORMANCE OF EIGCC
C
      TYPE 400,WK2(1)
      IF(WK2(1).GT.1.) TYPE 400,WK2(1)
400  FORMAT (' EIGCC PERFORMANCE=',1PE10.3)
C
C      FIND INVERSE OF U
C
      DO 410 I=1,NM
      DO 410 J=1,NM
410  V(I,J)=U(I,J)
      CALL LEQT2C(V,NM,MATSIZ,DUM,1,MATSIZ,1,WK2,IER)
      IF(IER.NE.129) GO TO 412
      TYPE 411
411  FORMAT (' U IS SINGULAR')
      STOP
412  IF(IER.EQ.130) TYPE 413
413  FORMAT (' ITERATION FAILED TO IMPROVE U INVERSE. MATRIX IS',
1' TOO ILL CONDITIONED')
      DO 414 IM=1,NM
      DO 415 JM=1,NM
415  DUM(JM,1)=0.
      DUM(IM,1)=1.
      CALL LEQT2C(V,NM,MATSIZ,DUM,1,MATSIZ,2,WK2,IER)
      IF(IER.NE.129) GO TO 416
      TYPE 411
      STOP
416  IF(IER.EQ.130) TYPE 413
      DO 417 JM=1,NM
417  UI(JM,IM)=DUM(JM,1)
414  CONTINUE

```

```

C      TYPE 217,((UI(I,J),J=1,NM),I=1,NM)
217 FORMAT(' 217 UI'/(6G))
C
C      FORM MATRIX SIGMA*S1
      DO 115 I=1,NM
      DO 115 J=1,NM
115 DUM(I,J)=SIG(I,I)*U(I,J)
C      TYPE 218,((DUM(I,J),J=1,NM),I=1,NM)
218 FORMAT(' 218 DUM'/(6G))
C
C      FORM MATRIX U=S*SIGMA*S1
C
      DO 117 I=1,NM
      DO 117 J=1,NM
      SUM=0.
      DO 118 K=1,NM
118 SUM=SUM+S(I,K)*DUM(K,J)
117 U(I,J)=SUM
C      TYPE 219,((U(I,J),J=1,NM),I=1,NM)
219 FORMAT(' 219 U'/(6G))
C
C      FORM MATRIX S1 INVERSE*SIGMA INVERSE
C
      DO 119 I=1,NM
      DO 119 J=1,NM
119 DUM(I,J)=UI(I,J)/SIG(J,J)
C      TYPE 220,((DUM(I,J),J=1,NM),I=1,NM)
220 FORMAT(' 220 DUM'/(6G))
C
C      FORM MATRIX UI INVERSE=S1 INVERSE*SIGMA INVERSE*S INVERSE
C
      DO 120 I=1,NM
      DO 120 J=1,NM
      SUM=0.
      DO 121 K=1,NM
121 SUM=SUM+DUM(I,K)*SI(K,J)
120 UI(I,J)=SUM
C      TYPE 221,((UI(I,J),J=1,NM),I=1,NM)
221 FORMAT(' 221 UI'/(6G))
      RETURN
      END

```

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TY COEF1

SUBROUTINE COEF1(NM,AM,V,A,AO,ALF,BET,FRAC)  
DIMENSION AM(30,30),V(15),A(15),AO(15),ALF(3,15),BET(3,15),  
IC(15,15),V1(15),V2(15),V3(15)  
COMPLEX AM,V,A,AO,ALF,BET,C,SUM,SUM2

ROUTINE COMPUTES COEFFICIENTS IN INFINITE AND SOURCE SEGMENTS BY  
MATCHING, USING AN ITERATIVE APPROACH.

NC=10

TYPE 911

911 FORMAT(' INPUT MAXIMUM NUMBER OF ITERATIONS')

ACCEPT 912,NC

912 FORMAT(G)

FOR MODES WHICH DO NOT PROPAGATE AT SOURCE, V=0.

SET V=.00001\*SMALLEST NONZERO V FOR ITERATION PROCESS.

VMIN=CABS(V(1))

DO 100 I=2,NM

VT=CABS(V(I))

IF(VT.LT.VMIN.AND.VT.GT.0.) VMIN=VT

100 CONTINUE

VMIN=VMIN\*1.E-5

DO 101 I=1,NM

IF(V(I).EQ.(0.,0.)) V(I)=VMIN/(0.,1.)

V1(I)=REAL(V(I))

V2(I)=AIMAG(V(I))

V3(I)=FRAC\*CABS(V(I))

101 CONTINUE

TYPE 229,(V(I),I=1,NM)

229 FORMAT(' 229 V'/60)

TYPE 230,(V1(I),I=1,NM)

230 FORMAT(' 230 V1'/30)

TYPE 231,(V2(I),I=1,NM)

231 FORMAT(' 231 V2'/30)

TYPE 232,(V3(I),I=1,NM)

232 FORMAT(' 232 V3'/30)

COMBINE ELEMENTS OF AM FOR USE IN EQUATION 4 OF PAGE NAGL.4

DO 102 I=1,NM

I1=2\*(I-1)

DO 102 J=1,NM

102 C(I,J)=AM(I1+1,J)-AM(I1+2,J)

TYPE 233,((C(I,J),J=1,NM),I=1,NM)

233 FORMAT(' 233 C'/(60))

DEFINE A FOR ZEROTH ORDER (PAGE NAGL.5)

DO 103 I=1,NM

103 A(I)=V(I)/C(I,I)

TYPE 234,(A(I),I=1,NM)

234 FORMAT(' 234 A'/60)

IC=ITERATION COUNTER

IC=1

GO TO 50

DEFINE A FOR IC-TH ITERATION

4 DO 104 I=1,NM

SUM=0.

DO 105 J=1,NM

IF(J,EQ,I) GO TO 107

SUM=SUM+C(I,J)\*A(J)  
105 CONTINUE  
104 A(I)=(V(I)-SUM)/C(I,I)

C  
C  
C

THE FOLLOWING CARD IS TEMPORARY TO TEST ITERATION PROCEDURE

50 IW=0  
DO 106 I=1,NM  
SUM=0.  
DO 107 J=1,NM  
107 SUM=SUM+C(I,J)\*A(J)  
X1=REAL(SUM)  
X2=AIMAG(SUM)

C  
C  
C

THE FOLLOWING 2 CARDS SHOULD BE INCLUDED WHEN ITERATION TEST  
IS COMPLETED

C  
C  
C

IF(ABS(X1-V1(I)).GT.V3(I)) GO TO 1  
IF(ABS(X2-V2(I)).GT.V3(I)) GO TO 1

C  
C  
C

THE FOLLOWING 3 CARDS ARE TEMPORARY TO TEST ITERATION PROCEDURE

C  
C  
C

IF(ABS(X1-V1(I)).GT.V3(I)) IW=1  
IF(ABS(X2-V2(I)).GT.V3(I)) IW=1  
TYPE 875,IC,I,A(I),X1,X2,V1(I),V2(I)  
875 FORMAT(2I3,1P2E15.7/6X,1P4E15.7)  
106 CONTINUE

C  
C  
C

THE FOLLOWING CARD SHOULD BE INCLUDED WHEN ITERATION TEST IS  
COMPLETED

C  
C  
C

GO TO 2

C  
C  
C

THE FOLLOWING CARD IS TEMPORARY TO TEST ITERATION PROCEDURE

C  
C  
C

IF(IW.EQ.0) GO TO 2  
1 IF(IC.EQ.NC) GO TO 3

C  
C  
C

START NEW ITERATION

C  
C  
C

IC=IC+1  
GO TO 4

C  
C  
C

MAXIMUM NUMBER OF ITERATIONS REACHED WITHOUT CONVERGING

3 TYPE 5  
5 FORMAT(// ' A NOT CONVERGED. SUM, V=' )  
DO 109 I=1,NM  
SUM=0.  
DO 110 J=1,NM  
110 SUM=SUM+C(I,J)\*A(J)  
TYPE 875,I,IC,A(I),SUM,V(I)  
109 CONTINUE  
TYPE 600  
600 FORMAT(' TYPE...-1 STOP//

1' 0 USE LAST VALUES FOR COEFFICIENTS//  
2' N ITERATE N MORE TIMES')  
ACCEPT 601,NC  
601 FORMAT(U)  
IF(NC) 602,2,603  
602 STOP  
603 IC=1  
GO TO 4

C  
C  
C

USING CONVERGED VALUES OF A (OR VALUES FROM NC-TH ITERATION),  
DEFINE AO,AND ALL DELTA FOR SAME INTERVAL



```
2 DO 111 I=1,NM
  I1=2*(I-1)
  SUM=0.
  SUM2=0.
  DO 112 J=1,NM
    SUM=SUM+AM(I1+1,J)*A(J)
112 SUM2=SUM2+AM(I1+2,J)*A(J)
  S=SQRT((CABS(SUM))**2+(CABS(SUM2))**2)
  A0(I)=S
  ALF(1,I)=SUM/S
111 BET(1,I)=SUM2/S
  TYPE 690,(A0(I),I=1,NM)
690 FORMAT(' 690 A0'/6G)
  TYPE 235,(ALF(1,I),I=1,NM)
235 FORMAT(' 235 ALF'/6G)
  TYPE 236,(BET(1,I),I=1,NM)
236 FORMAT(' 236 BET'/6G)
  RETURN
  END
```

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TY HMDIAG

SUBROUTINE HMDIAG(N,NP,NQ,CMAT,CEIG,CVEC,CVECC,DMAT,EMAT,DVEC,  
1DEIG,WK)  
COMPLEX CMAT,CEIG,CVEC,CVECC,EYE  
COMPLEX DMAT,EMAT,DVEC  
DIMENSION CMAT(15,15),CVEC(15,15),CVECC(15,15),CEIG(15)  
DIMENSION DMAT(N,N),EMAT(NP),DVEC(N,N),DEIG(N),WK(NQ)

C  
C  
C  
C  
C  
C

MUST DIVIDE BY EYE BECAUSE MATRIX IS ANTISYMMETRIC.  
THIS MAKES MATRIX HERMITIAN WHICH CAN BE EASILY  
DIAGONALIZED. ORIGINAL MATRIX HAS EIGENVALUES WITH SAME MAGNITUDE  
AS THE HERMITIAN MATRIX BUT ITS EIGENVALUES ARE PURELY IMAGINARY.  
THE EIGENVECTORS ARE THE SAME.

EYE=(0.,1.)  
DO 150 I=1,N  
DO 150 J=1,N  
150 DMAT(I,J)=CMAT(I,J)/EYE  
CALL VCVTCH(DMAT,N,N,EMAT)  
CALL EIGCH(EMAT,N,1,DEIG,DVEC,N,WK,IER)

C  
C  
C  
C

MUST MULTIPLY BY EYE BECAUSE ANTISYMMETRIC MATRIX HAS IMAGINARY  
EIGENVALUES

DO 151 I=1,N  
CEIG(I)=DEIG(I)\*EYE  
DO 151 J=1,N  
CVEC(I,J)=DVEC(I,J)  
151 CVECC(J,I)=CONJG(CVEC(I,J))  
RETURN  
END

TY MXEL

```

00050 SUBROUTINE MXEL(NMODE,NPRS,D,OXEL,OZB)
00100 IMPLICIT DOUBLE PRECISION (A-H,P-Z)
00200 DIMENSION X(15,4),A(15,4),B(15,4),U(15,4),UP(15,4),UNIT(
15)
00250 DIMENSION D(4),SL(4),AL(4),S(15,15),OXEL(15,15),OZB(4),N
MODE(4)
00250 M=NMODE(NPRS+NC-1)
00350 DO 36 NC=1,2
00360 DO 35 NP=1,NPRS+NC-1
00300 DO 30 J=1,M
00400 READ(21,40) NB
00500 DO 30 I=1,NB
00600 30 READ(21,10) X(J,I),A(J,I),B(J,I),U(J,I),UP(J,I)
00700 10 FORMAT(5D)
00800 40 FORMAT(I4)
00900 C TYPE 80,((U(J,I),I=1,4),J=1,M)
01000 80 FORMAT(4D16.7)
01100 C TYPE 80,((UP(J,I),I=1,4),J=1,M)
01200 35 CONTINUE
01300 IF(NC.EQ.2) GO TO 39
01400 READ(21,40) NB1
01500 DO 31 I=1,NB1
01600 31 READ(21,10) C1,AL(I),C2,C3,C4
01700 DO 38 KL=2,NB
01800 38 SL(KL)=(AL(KL)**3-A(1,KL)**3)/
1 (D(NPRS+1)-D(NPRS))
02000 39 CONTINUE
02100 CLOSE(UNIT=21)
02600 DO 100 J=2,M
02700 DO 100 K=1,J-1
02800 IF(NC.EQ.1) S(J,K)=0.
02900 DO 110 I=2,NB
03000 BD=(B(J,I)-B(K,I))*A(J,I)**2
03100 BS=B(J,I)+B(K,I)
03200 V=A(J,I)*(BS+2.*X(J,I))*U(J,I)*U(K,I)+2.*UP(J,I)*UP(K,I)
03300 W=(BD*X(J,I)-2.*A(J,I)/BD)*(UP(J,I)*U(K,I)-U(J,I)*UP(K,I)
))
03400 Y=A(J,I)*(BS+2.*X(J,I-1))*U(J,I-1)*U(K,I-1)+
1 2.*UP(J,I-1)*UP(K,I-1)*(A(J,I-1)/A(J,I))**2
03600 Z=(X(J,I-1)*BD-2.*A(J,I)/BD)*(UP(J,I-1)*U(K,I-1)-
1 U(J,I-1)*UP(K,I-1))*A(J,I-1)/A(J,I)
03800 S(J,K)=S(J,K)+(V+W-Y-Z)/BD/BD*SL(I)/2./(OZB(NPRS+NC-1))*
*2
03850 OXEL(J,K)=S(J,K)
03860 105 FORMAT(2I4,D)
03870 110 CONTINUE
03900 100 CONTINUE
03950 36 CONTINUE
04000 CONTINUE
02500 RETURN
02600 END

```

TY HANK

SUBROUTINE HANK(RHO,H1,H1P,H2,H2P)  
COMPLEX H1,H1P,H2,H2P  
CALL HANLG(RHO,H1,H1P,H2,H2P)  
RETURN  
END

TY HANLG

00100		SUBROUTINE HANLG(X,H1,H1P,H2,H2P)
00200		REAL J0,J1
00300		DOUBLE PRECISION PI,Z,A,P,Q,R,S,X
00400		COMPLEX H1,H1P,H2,H2P
00500		PI=3.14159265357
00600		Z=9*X
00700		Z1=DSQRT(2/PI/X)
00800		A=X-PI/4
00900		P=1-4.5/Z**2+9*25*49/24./Z**4
01000		Q=-1/Z+9*25/6./Z**3
01100		R=1+7.5/Z**2
01200		S=3/Z-9*35/6./Z**3
01300		J0=Z1*SNGL(P*DCOS(A)-Q*DSIN(A))
01400		Y0=Z1*SNGL(P*DSIN(A)+Q*DCOS(A))
01500		H1=J0+(0.,1.)*Y0
01600	50	FORMAT(2G)
01700		J0=-Z1*(R*DSIN(A)+S*DCOS(A))
01800		Y0=Z1*(R*DCOS(A)-S*DSIN(A))
01900		H1P=J0+(0.,1.)*Y0
02000		H2=CONJG(H1)
02100		H2P=CONJG(H1P)
02200		RETURN
02300		END



TY DAIRY

```

00010 SUBROUTINE DAIRY(DX,AI,AIP,BI,BIP)
26230
00020 IMPLICIT DOUBLE PRECISION (A-H,O-Z)
00030 C FOR DOUBLE PRECISION ARGUMENTS, THIS ROUTINE CALCULATES THE A
IRY 26240
00040 C FUNCTION AI(X) AND ITS DERIVATIVE AIP(X). IT ALSO FINDS
26250
00050 C THE OTHER REAL LINEARLY INDEPENDENT SOLUTION BI(X) AND
26260
00060 C ITS DERIVATIVE BIP(X).
26270
00070 C THE DEFINITIONS AND NORMALIZATIONS ARE AS IN NBS HANDBOOK
26280
00080 C OF MATHEMATICAL FUNCTIONS,P.446
26290
00090 C THE METHODS USED ARE POWER SERIES EXPANSION FOR SMALL X
26300
00100 C AND GAUSSIAN INTEGRATION FOR LARGE X
26310
00110 DIMENSION X(16),W(16),XSQ(16)
26320
00120 C DOUBLE PRECISION DX,AI,AIP,BI,BIP
26330
00130 C DOUBLE PRECISION XS,XCUBE,AISUM,AIPSUM
26340
00140 C DOUBLE PRECISION DF,DFP,DG,DGP
26350
00150 C DOUBLE PRECISION FJM2,FJM1,FJ,FJP1,FJP2,FACTOR
26360
00160 C DOUBLE PRECISION C1,C2,ROOT3
26370
00170 C DOUBLE PRECISION DZETA,DARG,DROOTX
26380
00180 C DOUBLE PRECISION ROOT4X,S,CD,RATIO,EFAC,ZETASQ
26390
00190 C DOUBLE PRECISION SUMR,SUMI,SUMRP,SUMIP,TERMR,TERMI
26400
00200 C DOUBLE PRECISION DZERO,DA,DB,DEN,ONE
26410
00210 C DOUBLE PRECISION X,W,XSQ
26420
00220 C DOUBLE PRECISION RSQ,TEMP, RTPI,RTPI2
26430
00230 C DOUBLE PRECISION TERMA,TERMB
26440
00240 LOGICAL NEEDDI
26450
00250 DATA DZERO,ONE /0.0D0,1.0D0/
26460
00260 DATA ROOT3/1.732050807568877D0/
26470
00270 DATA C1,C2 /.355028053887817D0, .258819403792807D0/
26480
00280 DATA RTPI /.2820947917738781D0/
26490
00290 DATA RTPI2/.5541895835477562D0/
26500
00300 C POSITIONS AND WEIGHTS FOR 10-TERM SUM FOR AIRY FUNCTIONS
26510
00310 DATA W( 1) / 3.154251762964787D-14/
26520
00320 DATA W( 2) / 6.6394219819584921D-11/
26530
00330 DATA W( 3) / 1.758381261355669D-08/
26540

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00340 DATA W( 4) / 1.3715592370435815D-06/  
26550  
00350 DATA W( 5) / 4.4350966632284350D-05/  
26560  
00360 DATA W( 6) / 7.1535010917718255D-04/  
26570  
00370 DATA W( 7) / 6.4889564103335381D-03/  
26580  
00380 DATA W( 8) / 3.6440415875773282D-02/  
26590  
00390 DATA W( 9) / 1.4397792418590999D-01/  
26600  
00400 DATA W(10) / 8.1231141336261486D-01/  
26610  
00410 DATA X( 1) / 1.4083081072180964D+01/  
26620  
00420 DATA X( 2) / 1.0214885479197331D+01/  
26630  
00430 DATA X( 3) / 7.4416018450450930D+00/  
26640  
00440 DATA X( 4) / 5.3070943061781927D+00/  
26650  
00450 DATA X( 5) / 3.6340135029132462D+00/  
26660  
00460 DATA X( 6) / 2.3310652303052450D+00/  
26670  
00470 DATA X( 7) / 1.3447970824609268D+00/  
26680  
00480 DATA X( 8) / 6.4188858369567296D-01/  
26690  
00490 DATA X( 9) / 2.0100345998121046D-01/  
26700  
00500 DATA X(10) / 8.0594359172052833D-03/  
26710  
00510 DATA XSQ( 1) /0.19833317248562170D 03/  
26720  
00520 DATA XSQ( 2) /0.10434388535311650D 03/  
26730  
00530 DATA XSQ( 3) /0.55377438020178170D 02/  
26740  
00540 DATA XSQ( 4) /0.28165249974668990D 02/  
26750  
00550 DATA XSQ( 5) /0.13206054139355800D 02/  
26760  
00560 DATA XSQ( 6) /0.54338651079380440D 01/  
26770  
00570 DATA XSQ( 7) /0.18084791929954200D 01/  
26780  
00580 DATA XSQ( 8) /0.41202095387883690D 00/  
26790  
00590 DATA XSQ( 9) /0.40402390924418070D-01/  
26800  
00600 DATA XSQ(10) /0.64954507303538390D-04/  
26810  
00610 C POSITIONS AND WEIGHTS FOR 4-TERM SUM FOR AIRY FUNCTIONS  
26820  
00620 DATA W(11) / 4.7763903057577263D-05/  
26830  
00630 DATA W(12) / 4.9914306432910959D-03/  
26840  
00640 DATA W(13) / 8.6167846973840312D-02/  
26850  
00650 DATA W(14) / 9.0879095845981102D-01/  
26860  
00660 DATA X(11) / 3.9198329554455091D+00/  
26870

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00670      DATA X(12) / 1.6915619004923504D+00/
26880
00680      DATA X(13) / 5.0275532467263018D-01/
26890
00690      DATA X(14) / 1.9247060562015692D-02/
26900
00700      DATA XSQ(11) /0.15365090398596670D 02/
26910
00710      DATA XSQ(12) /0.28613816631634610D 01/
26920
00720      DATA XSQ(13) /0.25276291648668180D 00/
26930
00730      DATA XSQ(14) /0.37044934027789980D-03/
26940
00740      C POSITIONS AND WEIGHTS FOR 2-TERM SUM FOR AIRY FUNCTIONS
26950
00750      DATA W(15) / 9.6807280595773604D-01/
26960
00760      DATA W(16) / 3.1927194042263958D-02/
26970
00770      DATA X(15) / 3.6800601866153044D-02/
26980
00780      DATA X(16) / 1.0592469382112375D+00/
26990
00790      DATA XSQ(15) /0.13542842977111070D-02/
27000
00800      DATA XSQ(16) /0.11220040761098810D 01/
27010
00810      IF(DX.LT.-5.0D0) GO TO 100
27020
00820      NEEDBI=.FALSE.
27030
00830      IF(DX.GT.3.7D0) GO TO 200
27040
00840      C THIS ROUTE FOR SMALLX, USING POWER SERIES.
27050
00850      C INITIALIZE
27060
00860      10 XS = DX*DX
27070
00870      XCUBE = XS *DX
27080
00880      XS = XS *0.5D0
27090
00890      DF = C1
27100
00900      DFP = C1*XS
27110
00910      DG = C2*DX
27120
00920      DGP = C2
27130
00930      AISUM = DF - DG
27140
00940      AIPSUM = DFP - DGP
27150
00950      BI = DF + DG
27160
00960      BIP = DFP + DGP
27170
00970      FJM2=-2.0D0
27180
00980      20 FJM2=FJM2+3.0D0
27190
00990      FJM1=FJM2+ONE
27200

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```

01000      FJ=FJM1+ONE
          27210
01010      FJP1=FJ+ONE
          27220
01020      FJP2=FJP1+ONE
          27230
01030      RATIO = XCUBE/FJ
          27240
01040      DF = DF*RATIO/FJM1
          27250
01050      DFP = DFP*RATIO/FJP2
          27260
01060      DG = DG*RATIO/FJP1
          27270
01070      DGP = DGP*RATIO/FJM2
          27280
01080      BI = BI + (DF+DG)
          27290
01090      BIP = BIP + (DFP+DGP)
          27300
01100      IF(NEEDBI) GO TO 80
          27310
01110      AISUM = AISUM + (DF-DG)
          27320
01120      AIPSUM = AIPSUM + (DFP-DGP)
          27330
01130      C CONVERGENCE TEST
          27340
01140      80 IF(DABS(DF).GT.1.0D-16) GO TO 20
          27350
01150      C CONVERGENCE. COMPUTE FUNCTIONS
          27360
01160      99 BI = ROOT3*BI
          27370
01170      BIP = ROOT3*BIP
          27380
01180      C THIS RETURNS IF X IS BETWEEN 3.7 AND 8.0, SINCE IN SUCH CASES
          27390
01190      MORE
          27400
01200      C ACCURATE VALUES OF AI AND AIP HAVE ALREADY BEEN FOUND BY GAUS
          27410
01210      C INTEGRATION
          27420
01220      IF(NEEDBI)RETURN
          27430
01230      AI = AISUM
          27440
01240      AIP = AIPSUM
          27450
01250      C RETURN
          27460
01260      100 DROOTX = DSQRT(-DX)
          27470
01270      ROOT4X = DSQRT(DROOTX)
          27480
01280      DZETA = -.6666666666666667*DX*DROOTX
          27490
01290      DARG = DZETA - .7853981633974483
          27500
01300      SUMR = DZERO
          27510
01310      SUMI = DZERO
          27520
01320      SUMRP = DZERO
          27530

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01330      SUMIP = DZERO
          27540
01340      C TEST TO SEE HOW MANY TERMS ARE NEEDED IN GAUSSIAN INTEGRATION
          27550
01350      IF(DX.LT.(-200.D0)) GO TO 140
          27560
01360      IF(DX.LT.(-15.D0)) GO TO 130
          27570
01370      C THIS CASE FOR DX BETWEEN -5.0 AND -15.0
          27580
01380      LIMLO=1
          27590
01390      LIMHI=10
          27600
01400      GO TO 149
          27610
01410      C THIS CASE FOR DX BETWEEN -15.0 AND -200.
          27620
01420      130 LIMLO=11
          27630
01430      LIMHI=14
          27640
01440      GO TO 149
          27650
01450      C THIS CASE FOR DX.LT.-200.
          27660
01460      140 LIMLO=15
          27670
01470      LIMHI=16
          27680
01480      149 ZETASQ=DZETA**2
          27690
01490      DO 150 K=LIMLO,LIMHI
          27700
          TERMR=W(K)/((ZETASQ+XSQ(K))**2)
          27710
01510      SUMR = SUMR + TERMR
          27720
01520      TERMR=TERMR*X(K)
          27730
01530      SUMI=SUMI+TERMR
          27740
01540      TERMR=TERMR*X(K)
          27750
01550      SUMRP=SUMRP+TERMR
          27760
01560      150 SUMIP=SUMIP+TERMR*X(K)
          27770
01570      SUMR=(SUMR*ZETASQ+SUMRP)*ZETASQ
          27780
01580      TEMP=SUMI*ZETASQ
          27790
01590      SUMI=(TEMP+SUMIP)*DZETA
          27800
01600      SUMRP=SUMRP*DZETA
          27810
01610      SUMIP=SUMIP-TEMP
          27820
01620      C FORM AIRY FUNCTIONS
          27830
01630      196 S = DSIN(DARG)
          27840
01640      CO = DCOS(DARG)
          27850
01650      RATIO = RTPI2/ROOT4X
          27860

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```

01650      AI = RATIO*(CO*SUMR + S*SUMI)
27870
01670      BI = RATIO*(CO*SUMI - S*SUMR)
27880
01680      SUMRP=SUMRP+SUMRP
27890
01690      RATIO = -.25D0/DX
27900
01700      FACTOR = -RTPI2*ROOT4X
27910
01710      AIP = RATIO*AI - DROOTX*BI + FACTOR*(CO*SUMRP+S*SUMIP)
27920
01720      BIP = RATIO*BI + DROOTX*AI + FACTOR*(CO*SUMIP-S*SUMRP)
27930
01730      RETURN
27940
01740      C GAUSSIAN INTEGRATION FOR LARGE POSITIVE X
27950
01750      200 DROOTX = DSQRT(DX)
27960
01760      DZETA = .6666666666666667*DX*DROOTX
27970
01770      EFAC = DEXP(-DZETA)
27980
01780      ROOT4X = DSQRT(DROOTX)
27990
01790      AI = DZERO
28000
01800      BI = DZERO
28010
01810      AIP = DZERO
28020
01820      BIP = DZERO
28030
01830      IF(DX.LT.8.0D0) NEEDBI=.TRUE.
28040
01840      C TEST TO SEE HOW MANY TERMS ARE NEEDED IN GAUSSIAN INTEGRATION
28050
01850      IF(DX.GT.15.0D0) GO TO 230
28060
01860      C THIS CASE FOR DX BETWEEN 3.7 AND 15.
28070
01870      LIMLO=1
28080
01880      LIMHI=10
28090
01890      GO TO 249
28100
01900      C THIS CASE FOR DX GREATER THAN 15.
28110
01910      230 LIMLO=11
28120
01920      LIMHI=14
28130
01930      249 DO 250 K=LIMLO,LIMHI
28140
01940      DA=DZETA+X(K)
28150
01950      TERMA = W(K)/DA
28160
01960      AI = AI + TERMA
28170
01970      AIP=AIP+TERMA*X(K)/DA
28180
01980      IF(NeedBI) GO TO 250
28190

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```

01990      DB=DZETA-X(K)
28200
02000      TERM = W(K)/DB
28210
02010      BI = BI + TERM
28220
02020      BIP=BIP+TERM*X(K)/DB
28230
02030  250  CONTINUE
28240
02040  C  FORM FUNCTIONS
28250
02050      FACTOR=RTPI*DZETA/ROOT4X
28260
02060      RATIO = 0.25D0/DX
28270
02070      AI=AI*EFAC*FACTOR
28280
02080      AIP=-(DROOTX+RATIO)*AI+RTPI*ROOT4X*EFAC*AIP
28290
02090  C  THIS IS SATISFIED ONLY FOR X BETWEEN 3.7 AND 9.0  IN THESE CA
SES
28300
02100  C  THE BI AND BIP ABOUT TO BE COMPUTED ARE NOT SUFFICIENTLY ACCU
RATE.
28310
02110  C  THUS RETURN TO POWER SERIES FOR BI AND BIP.
28320
02120      IF(NEEDEDI) GO TO 10
28330
02130      FACTOR=FACTOR+FACTOR
28340
02140      BI=BI*FACTOR/EFAC
28350
02150      BIP=(DROOTX-RATIO)*BI-RTPI2*ROOT4X*BIP/EFAC
28360
02160      RETURN
28370
02170      END
28380

```

```

TY FLT
  DIMENSION Y(150),IX(0/1000),IY(0/1000)
  TYPE 100
100 FORMAT(' PLOT TYPE=-1...NO MORE PLOTS'/
  114X,'0...TRANSMISSION LOSS'/
  212X,'N,1...N=MODE NUMBER, I=0(MAGNITUDE) OR 1(PHASE)')
  1 TYPE 101
101 FORMAT(' ENTER NUMBER OF MODES ON TAPE, PLOT TYPE')
  ACCEPT 102,NM,IPL,IMP
102 FORMAT(10G)
  NM=2*NM
  IF(IPL.EQ.-1) STOP
  TYPE 300
300 FORMAT(' TYPE TAPE NUMBER')
  ACCEPT 301,NT
301 FORMAT(G)
  IN=2*(IPL-1)+IMP+1
  TYPE 103
103 FORMAT(' ENTER 0 FOR LINE, 1 FOR POINTS')
  ACCEPT 102,ILPT
  TYPE 104
104 FORMAT(' INPUT NTOT,NMIN,NMAX,XMIN,XMAX,YMIN,YMAX,NX,NY')
  ACCEPT 102,NTOT,NMIN,NMAX,XMIN,XMAX,YMIN,YMAX,NX,NY
  IF(IPL.NE.0) GO TO 500
  YTEM=YMIN
  YMIN=-YMAX
  YMAX=-YTEM
500 CONTINUE
  DX=9999./NX
  DY=9999./NY
  TYPE 130
130 FORMAT(' PLTP')
  DO 10 I=0,NX
  IY(I)=0
  IX(I)=DX*I
  10 TYPE 30,IX(I),IY(I)
  30 FORMAT(2I10)
  DO 11 I=0,NY
  IX(I)=0
  IY(I)=DY*I
  11 TYPE 30,IX(I),IY(I)
  TYPE 70
  70 FORMAT(' PLTT')
  REWIND NT
  DX=9999./(XMAX-XMIN)
  DY=9999./(YMAX-YMIN)
  IF(IPL.NE.0) GO TO 200
  DO 12 I=1,NMIN
  READ (NT,105) X,Y(1)
105 FORMAT(F13.5,F15.6)
  IX(I)=DX*(X-XMIN)
  Y(1)=-Y(1)
  12 IY(1)=DY*(Y(1)-YMIN)
  GO TO 201
200 DO 202 I=1,NMIN
  READ (NT,203) X,(Y(J),J=1,NM)
203 FORMAT(F10.3,1P6E10.3/(10X,1P6E10.3))
  IX(I)=DX*(X-XMIN)
  202 IY(1)=DY*(Y(IN)-YMIN)
201 TYPE 130
  TYPE 30,IX(NMIN),IY(NMIN)
  TYPE 70
  IF(IPL.NE.0) GO TO 204
  DO 13 I=NMIN+1,NMAX
  READ (NT,105) X,Y(1)
  IX(I)=DX*(X-XMIN)
  
```



```
Y(1)=-Y(1)
13 IY(I)=DY*(Y(1)-YMIN)
GO TO 205
204 DO 206 I=NMIN+1,NMAX
READ (NT,203) X,(Y(J),J=1,NM)
IX(I)=DX*(X-XMIN)
206 IY(I)=DY*(Y(IN)-YMIN)
205 IF(ILFT.EQ.1) GO TO 222
TYPE 60
60 FORMAT(' PLTL')
GO TO 223
222 TYPE 130
223 CONTINUE
DO 14 I=NMIN,NMAX
14 TYPE 30,IX(I),IY(I)
TYPE 70
GO TO 1
END
```

TO RUN DEPTH PROGRAM.  
?TO?

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.EX ZAR25,DAIRY,SELAUT,PARAM,SERV,SERD,SINCOS

input: type 38  
output: 21,25,27,28  
also user 24,26,29

~C

.TO RUN RANGE PROGRAM:  
?TO?

.LOAD/F10 GPC,NAGL,FORM,COEF1,HMDIAG,MXEL,HANK,MANLG,DAIRY,STA:IMSL/LIB,  
STA:IMSL/LIB,STA:IMSL/LIB,STA:IMSL/LIB

input 21,25,27,28  
output 22,27

~C

.TO RUN PLOT PROGRAM:  
?TO?

.LOAD/F10 PLT

input 22,23

~C